

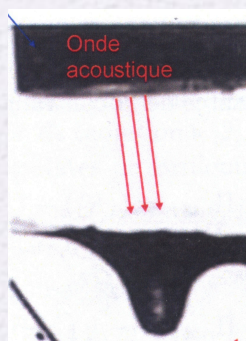
Acoustofluidics

Philippe Marmottant

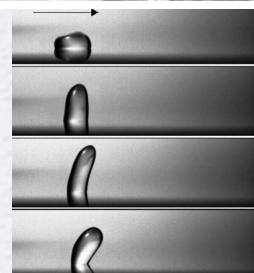
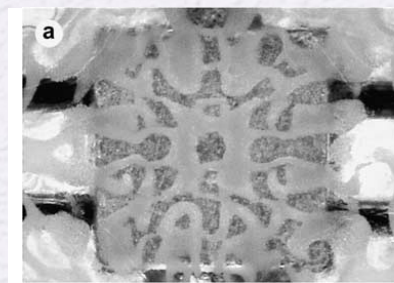
David Rabaud, Pierre Thibault, Q. Tran, Flore Mekki-Berrada, Thomas Combriat, Nicolas Bertin, Jean-François Louf, Maxime Harazi

Laboratoire Interdisciplinaire de
Physique (LIPhy) Grenoble

1) Using radiation pressure



2) Using acoustic streaming

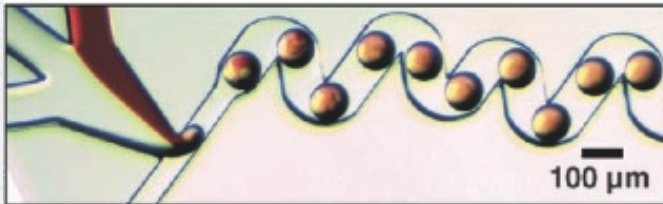


Microfluidics school Sète
October 13-18, 2019



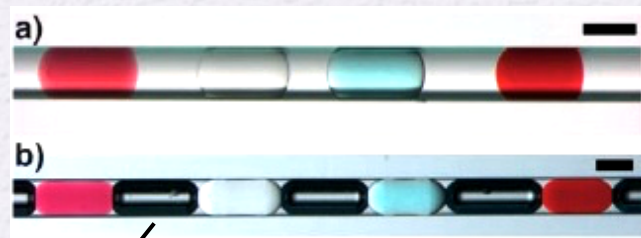
Context: digital microfluidics

Diluted droplets / bubbles follow the flow ...



Ismagilov group 2003

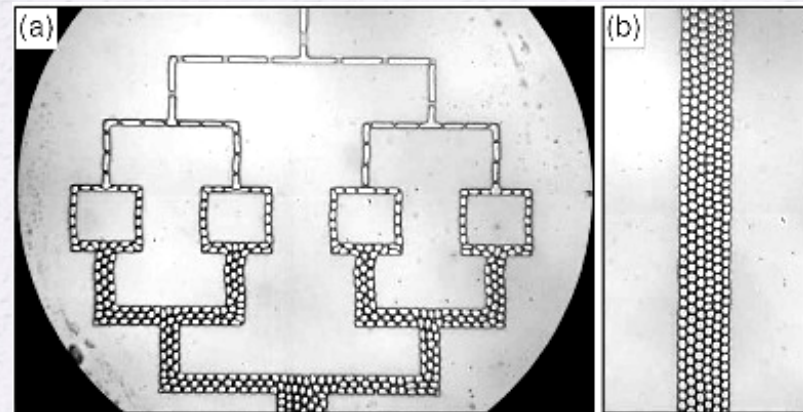
... or induce traffic jams



Bubble

Droplet

Ismagilov group 2005



Link 2004

How to use acoustic forces to control objects without contact?

Waves...

Acoustic wave definition: small disturbance around equilibrium of

- Pressure, $p = p_0 + p_1(x, t)$
- Velocity, $\mathbf{v} = \mathbf{v}_0 + \mathbf{v}_1(x, t)$

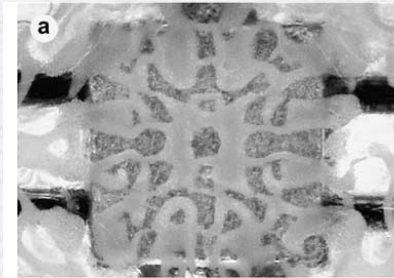
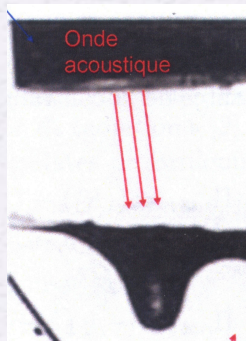
Propagating at speed of sound: $c = (K/\rho)^{1/2}$

with fluid compressibility $K = \rho(\partial p / \partial \rho)_{S=ct}$

How can sound exert forces and motion?

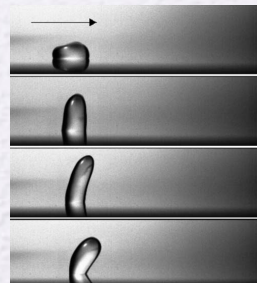
Outline

1) Radiation pressure



- a) on interfaces
- b) on particles, tweezers
- c) on bubbles

2) Acoustic streaming



- a) Quartz wind
- b) near boundaries
- c) near bubbles

1) Radiation pressure

- **Exists** in acoustics and in optics
(laser forces, as in optical tweezers)
- **Peculiarity in acoustics**
radiation pressure can be negative

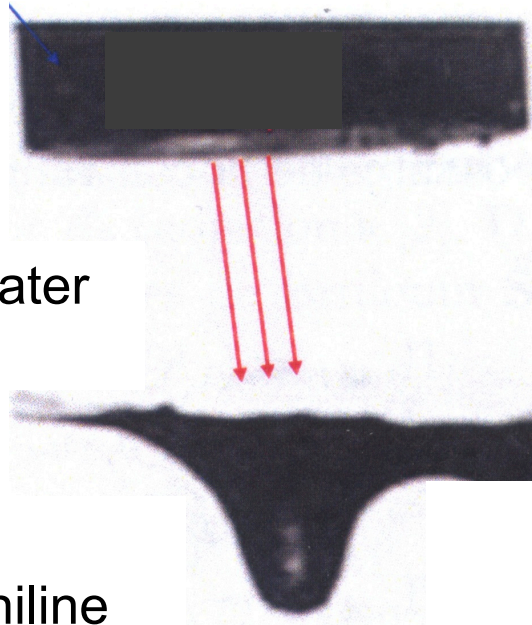
a) On fluid interfaces

Transducer

Water

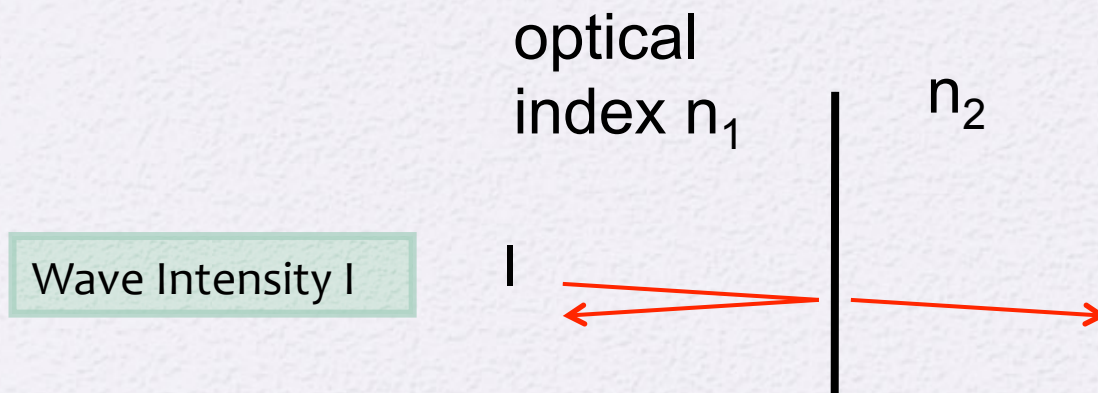
Aniline

Hertz and Mende 1939



Origin: change of momentum at reflection

- Case of optics

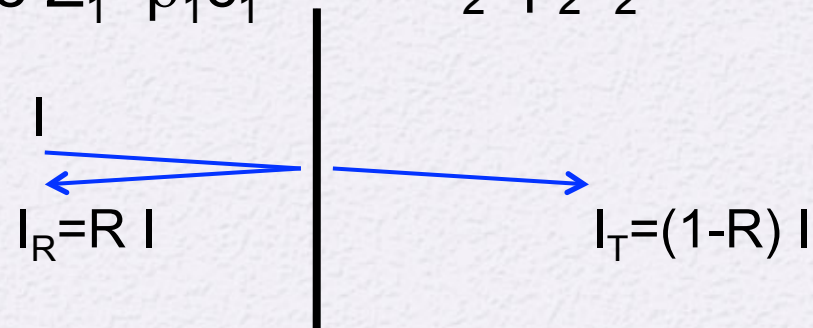


- Acoustics

acoustic

impedance $Z_1 = \rho_1 c_1$

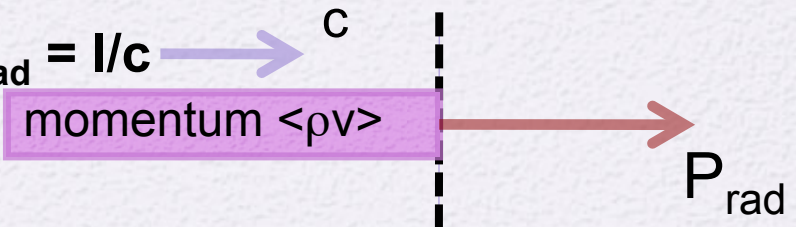
$Z_2 = \rho_2 c_2$



$$\text{Reflection } R = \frac{(Z_1 - Z_2)^2}{(Z_1 + Z_2)^2}$$

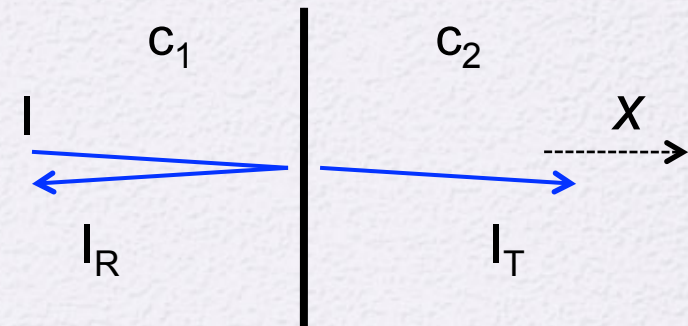
Wave Intensity is $I = P_A^2 / 2Z$ (W/m²)
with peak acoustic pressure P_A

Wave momentum

Flux of momentum results in a pressure $\mathbf{P}_{\text{rad}} = I/c$ 

Reflection on an interface

- incident wave: $\mathbf{P}_{\text{rad}}^{\text{incident}} = I/c_1$
- reflected wave: $\mathbf{P}_{\text{rad}}^{\text{reflected}} = I_R/c_1$
- transmitted wave: $\mathbf{P}_{\text{rad}}^{\text{transmitted}} = -I_T/c_2$

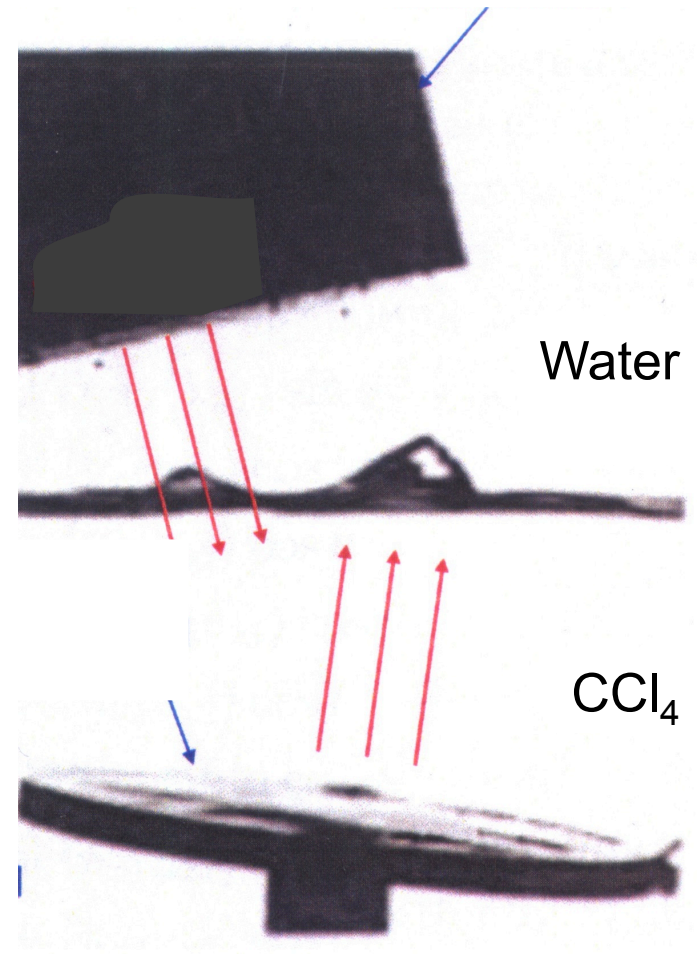
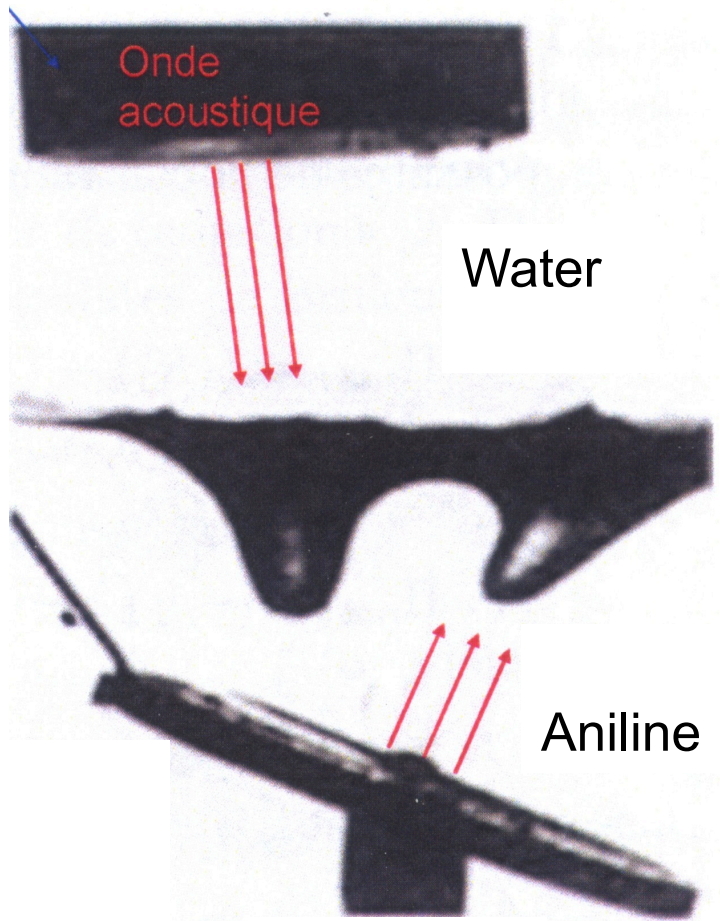


Total
along +x

$$\mathbf{P}_{\text{rad}} = I \left(\frac{1+R}{c_1} - \frac{1-R}{c_2} \right)$$

- \mathbf{P}_{rad} can be negative (e.g. if $c_2 \ll c_1$)
- Can be independent of the direction of propagation if exchanging $c_2 \leftrightarrow c_1$ changes sign of \mathbf{P}_{rad}

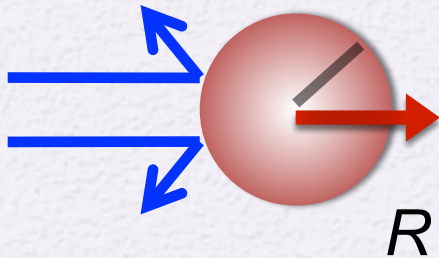
Transducer



Hertz and Mende 1939

b) On solid particles: acousto-phoresis

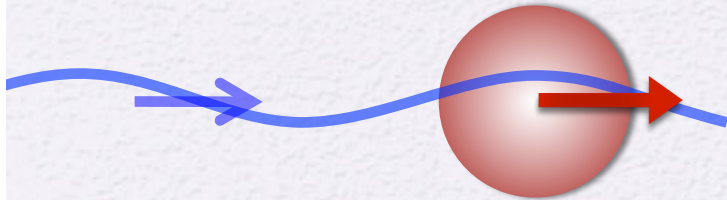
- Mie regime, $R \gg \lambda$



$$F_{\text{rad}} = S I/c$$

$$\text{with } S = \pi R^2$$

- Rayleigh regime, $R \ll \lambda$ (microfluidics)



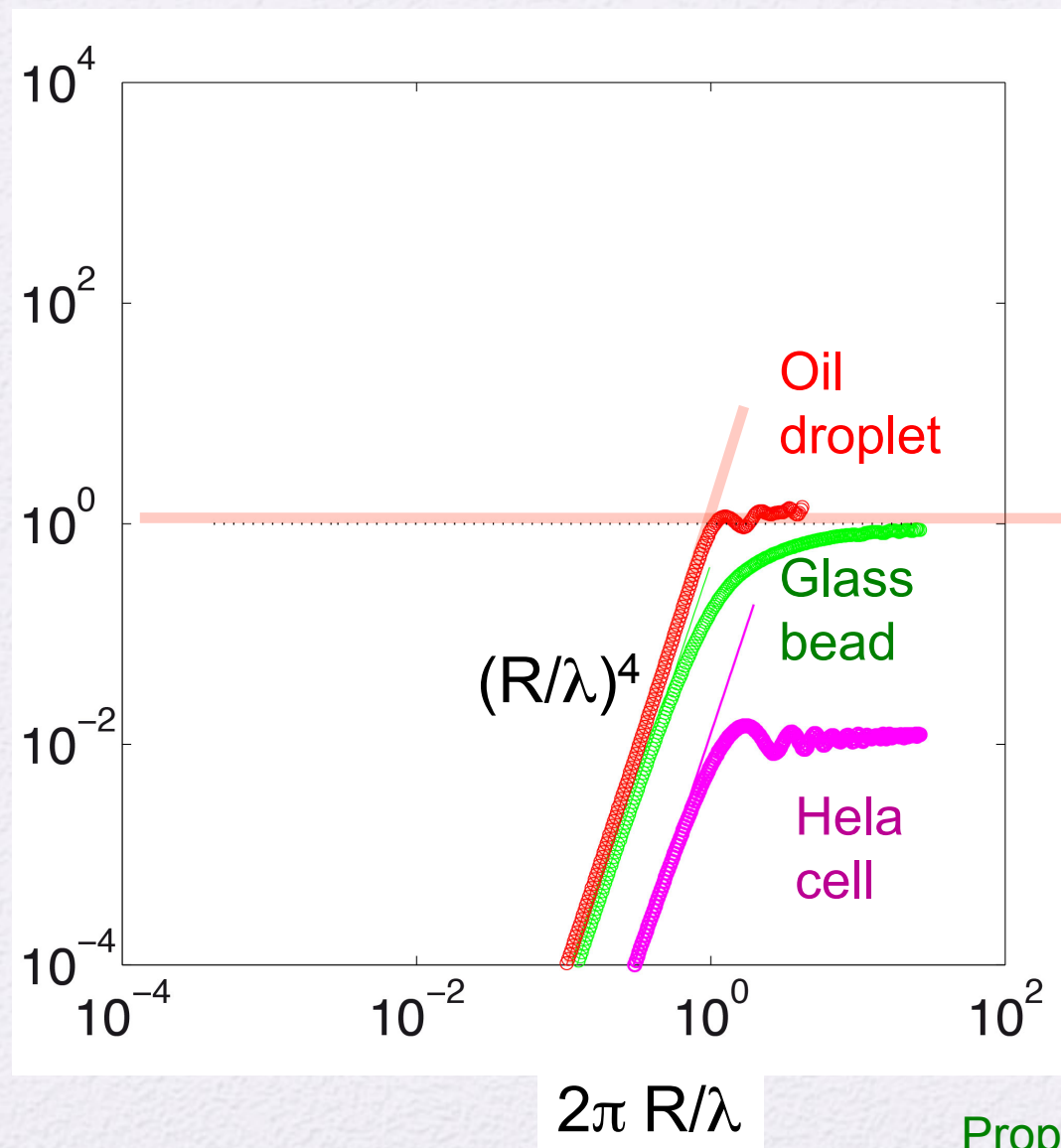
$$F_{\text{rad}} = \Phi_{\text{prop}} (R/\lambda)^4 S I/c$$

King 1934

Φ_{prop} depends on the particle properties

Yosioka 1955

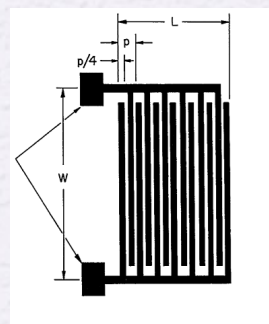
Radiation
force
 $F_{\text{rad}} / (\pi R^2 I/c)$



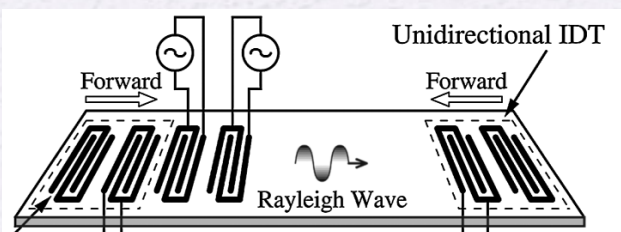
Propagating waves
Yosioka 1955

In practice:

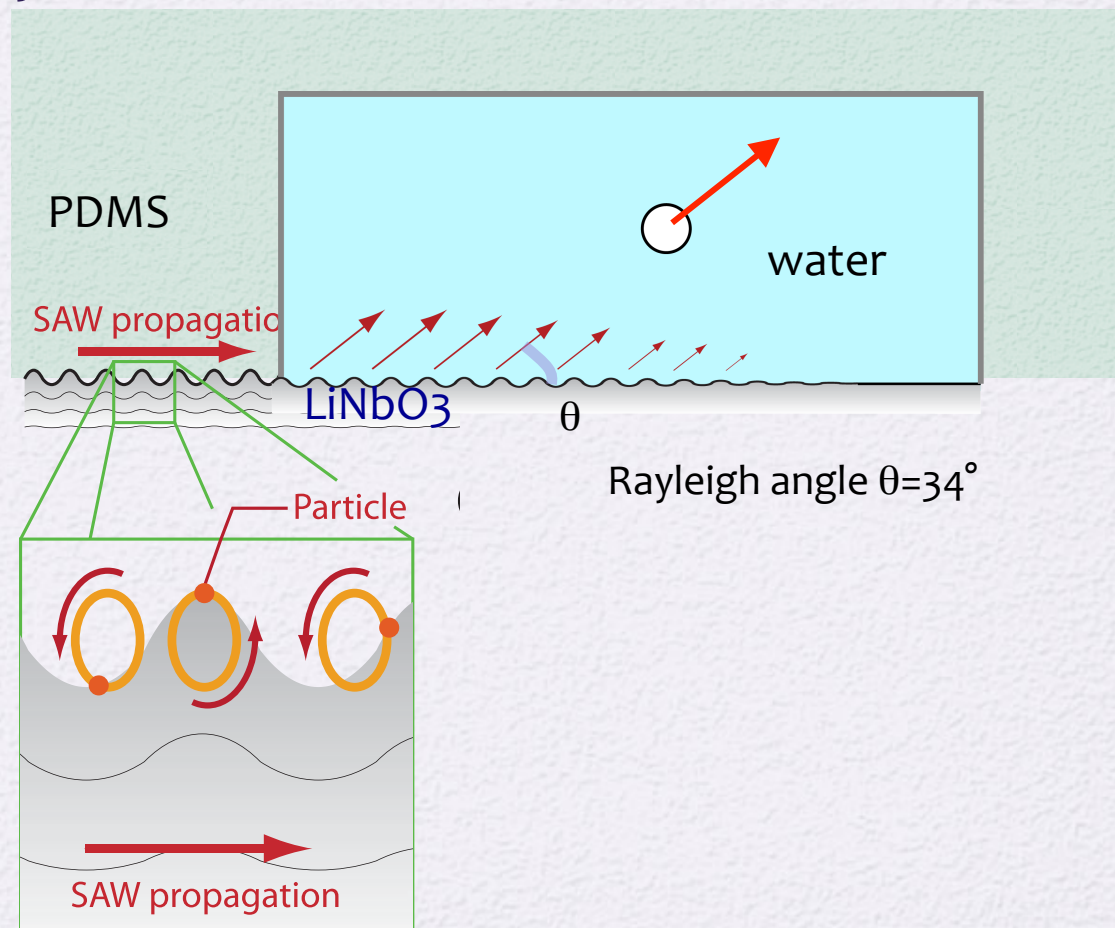
inject sound in a microsystem with surface acoustic waves



Interdigitated electrodes (IDT)



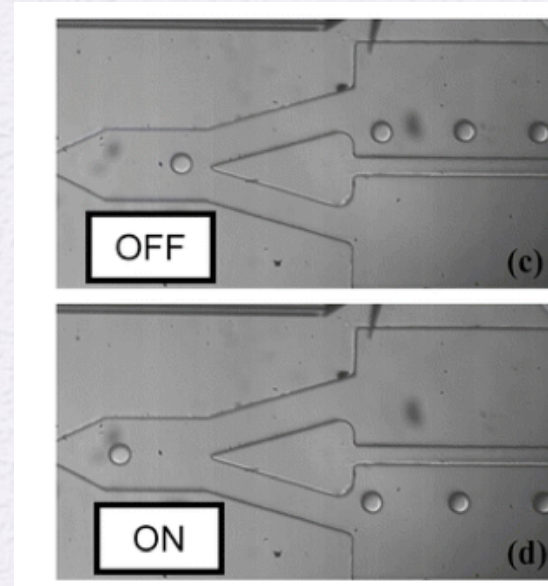
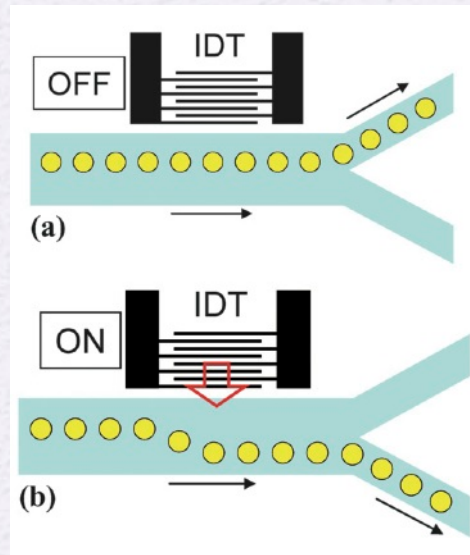
LiNbO_3
piezo-electric crystal
transparent



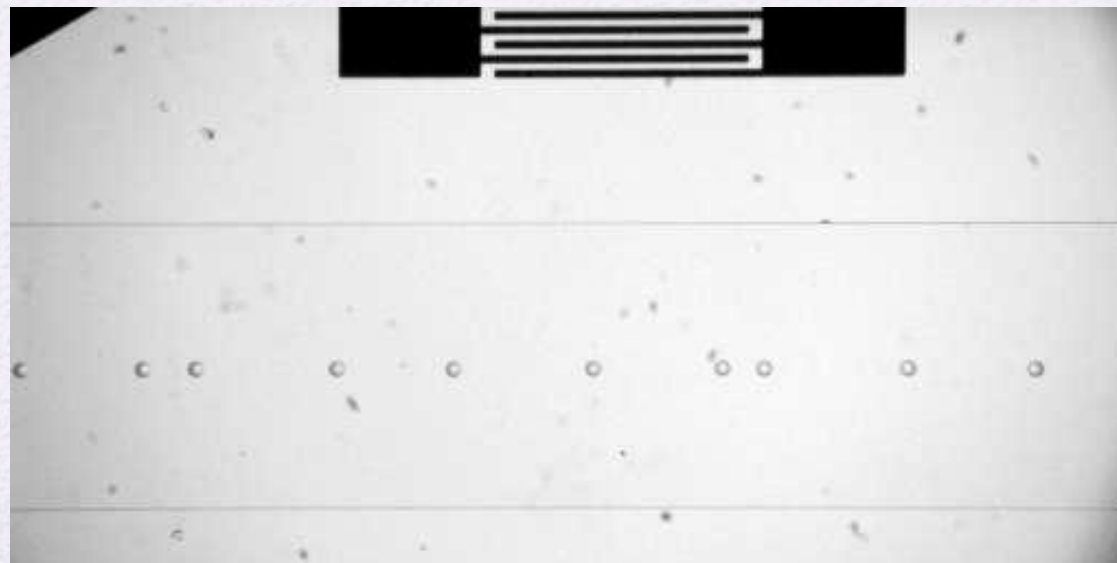
(SAW or Rayleigh waves)

See review by *Friend and Yeo 2011*

a tool to push particles/droplets



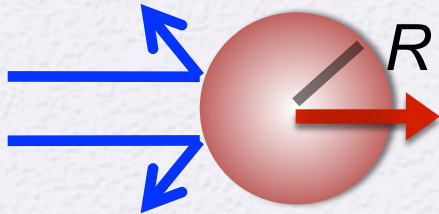
*Franke 2009
Lab on a Chip*



*our lab
unpublished*

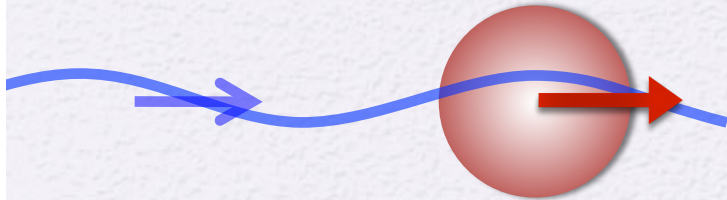
standing waves/propagating waves

- Mie regime, $R \gg \lambda$

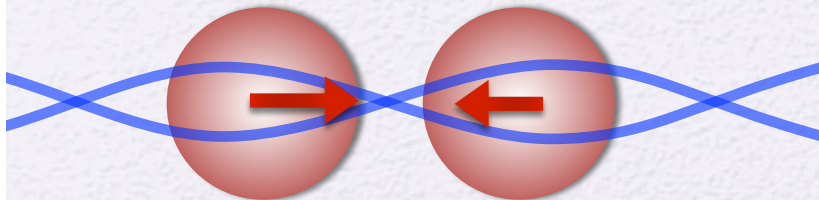


Propagative: $F_{\text{rad}} = S I/c$

- Rayleigh regime, $R \ll \lambda$ (microfluidics)

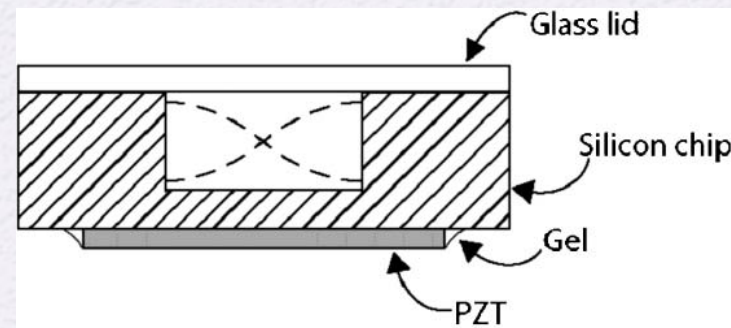
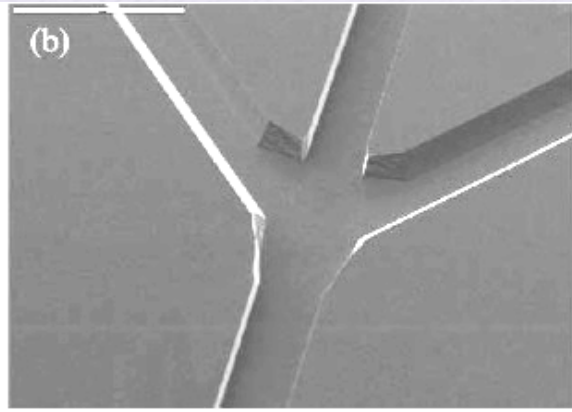


Propagative: $F_{\text{rad}} \sim (R/\lambda)^4 S I/c$



Stationary: $F_{\text{rad}} \sim (R/\lambda)^1 S I/c$

standing waves, examples

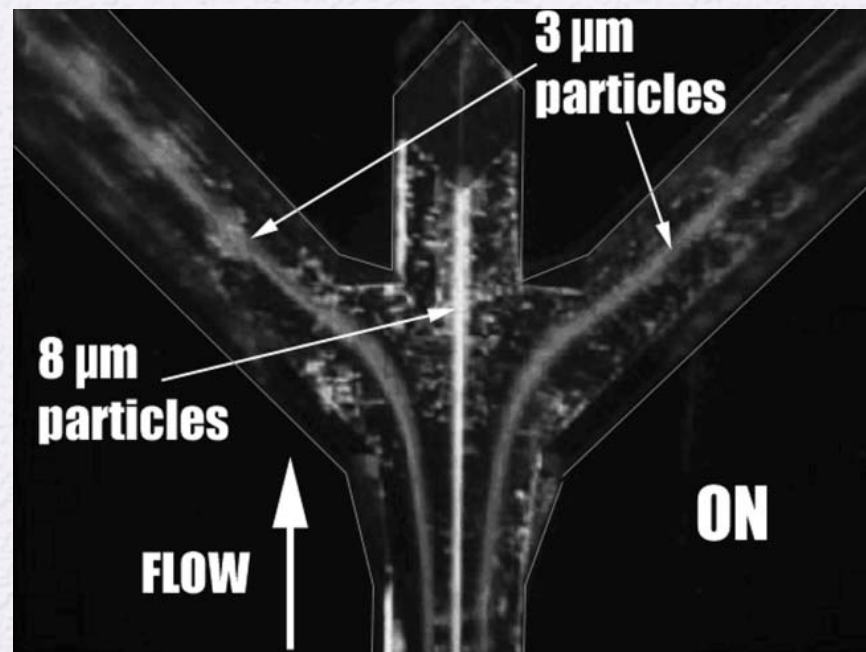


$$F_{\text{rad}} \sim (R/\lambda)^1 S I/c \sim R^3$$

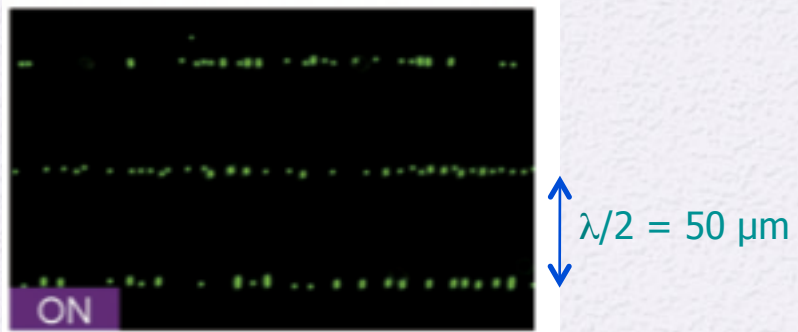
$$F_{\text{drag}} = 6\pi\eta Rv \sim R^1$$

stronger effect on large particles

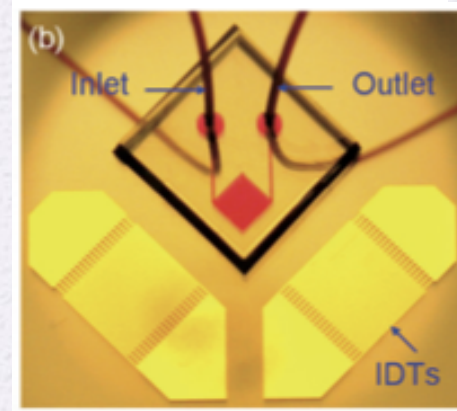
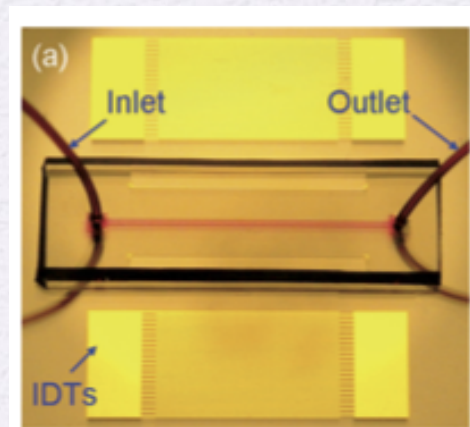
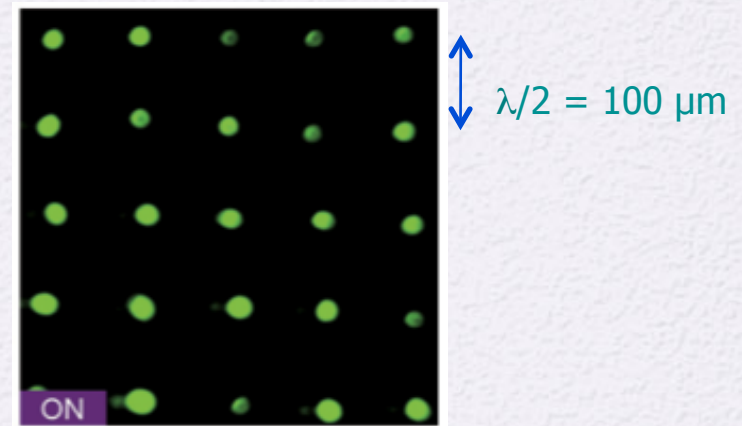
Petersonn 2007



1D - patterning



2D - patterning



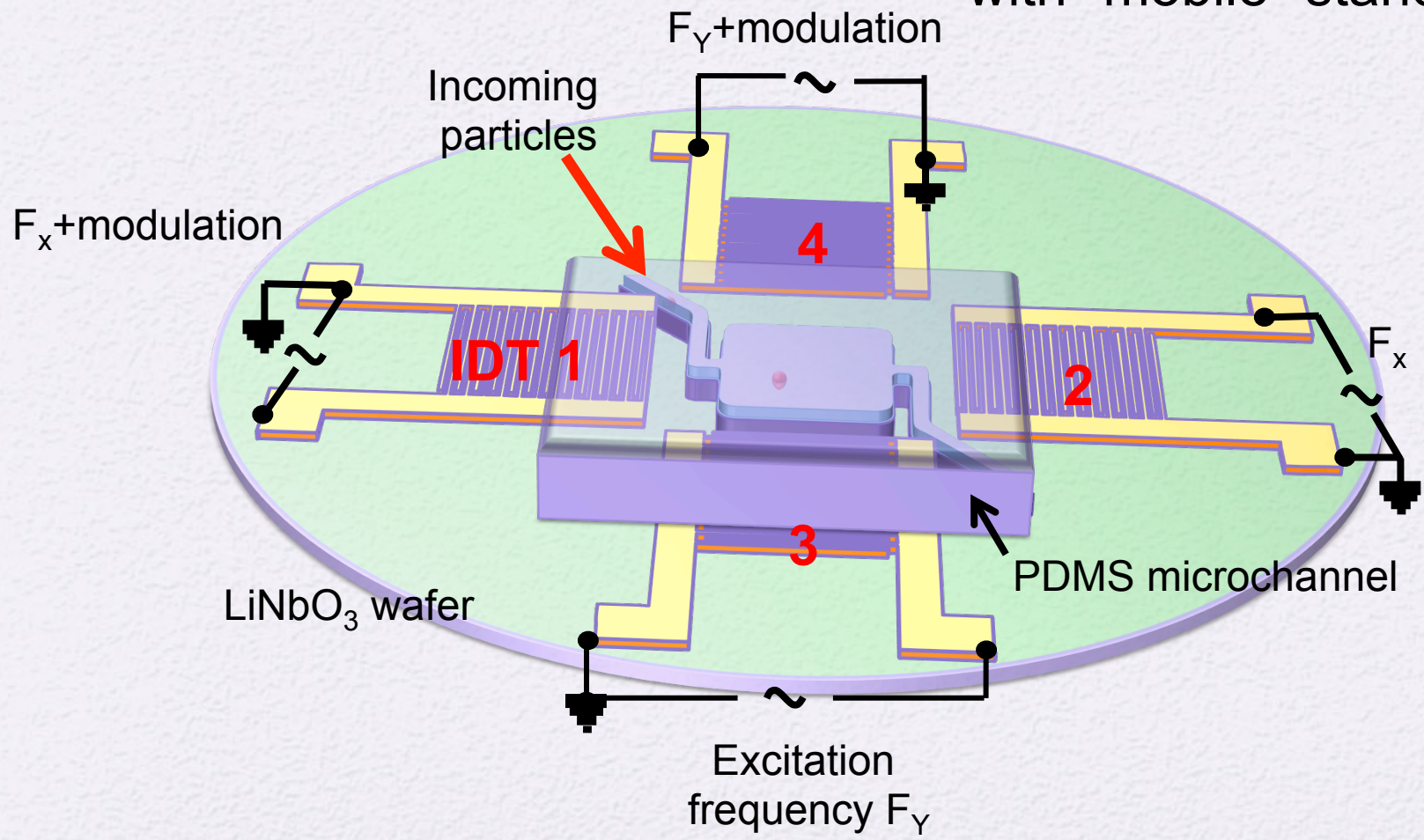
Shi 2009

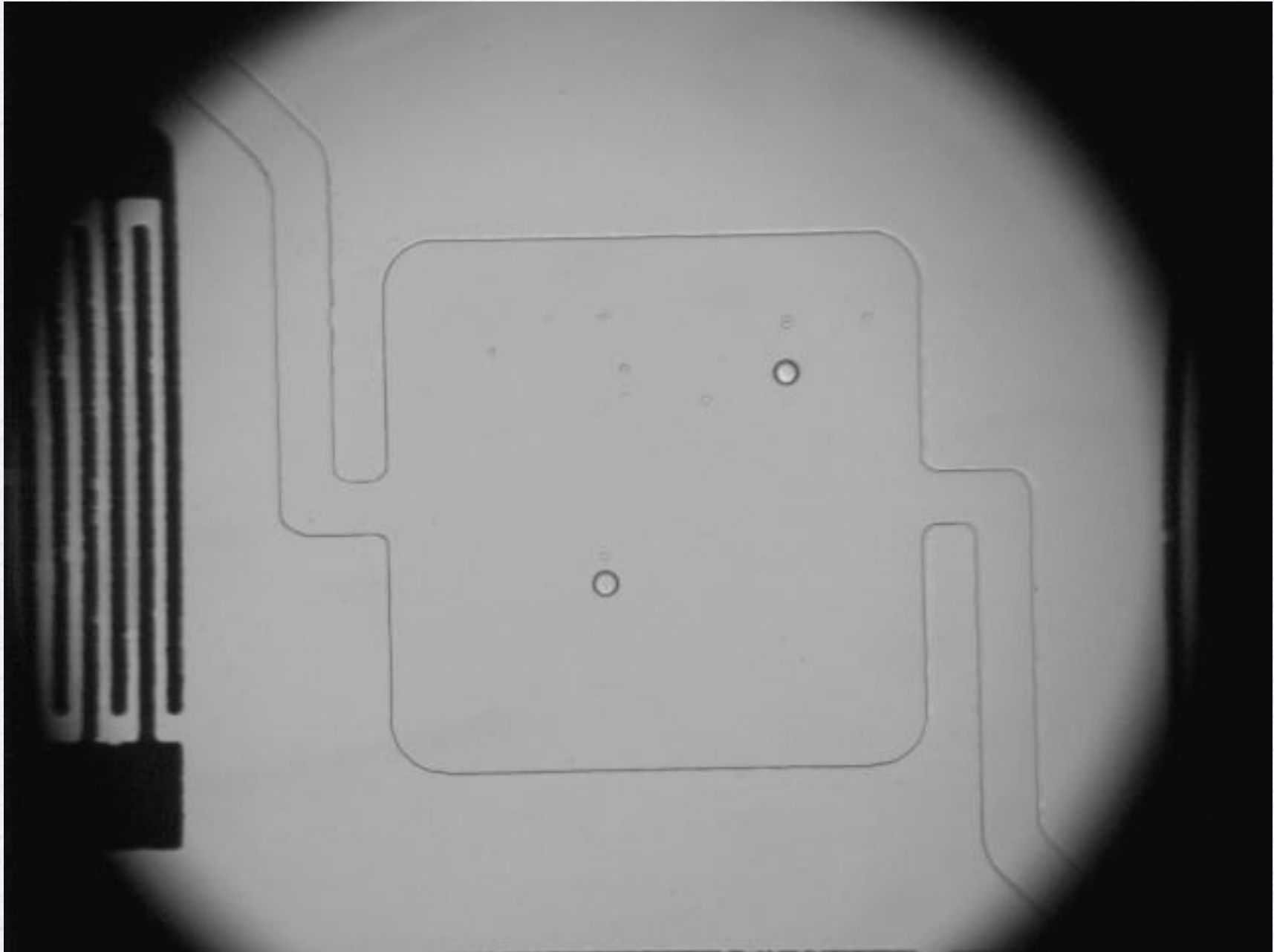
Acoustic sorting
ESPCI startup



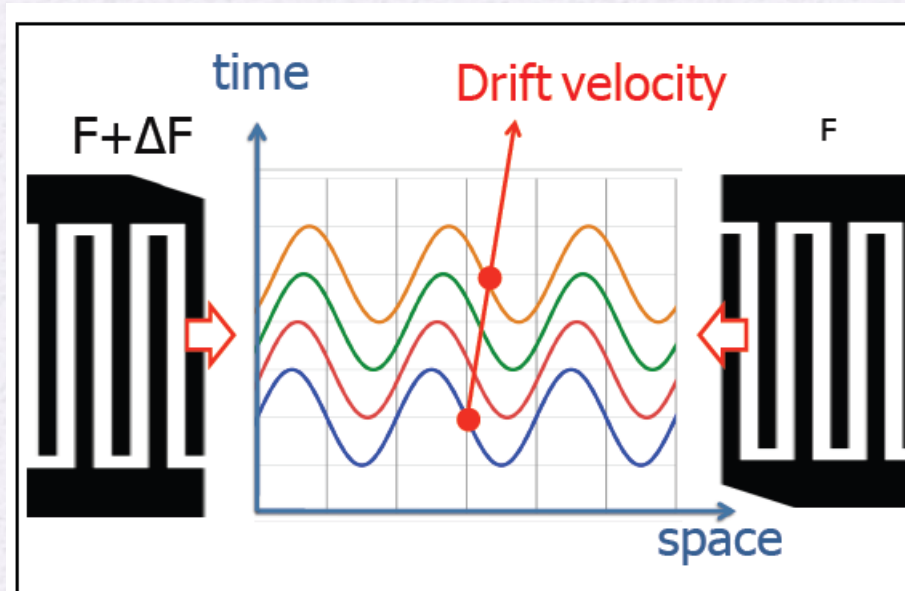
acoustic tweezers:

with "mobile" standing waves





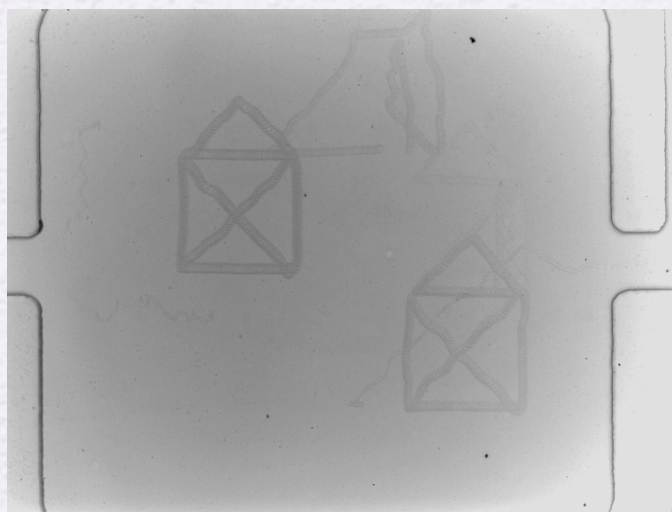
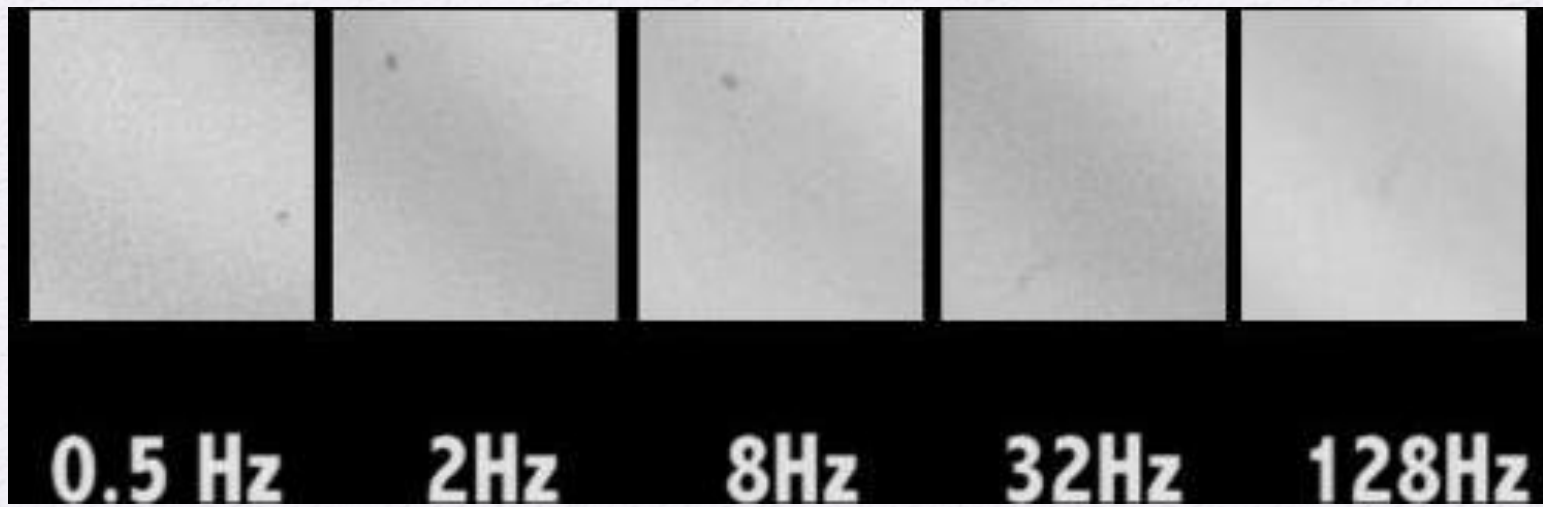
Principle: phase shift on one electrode



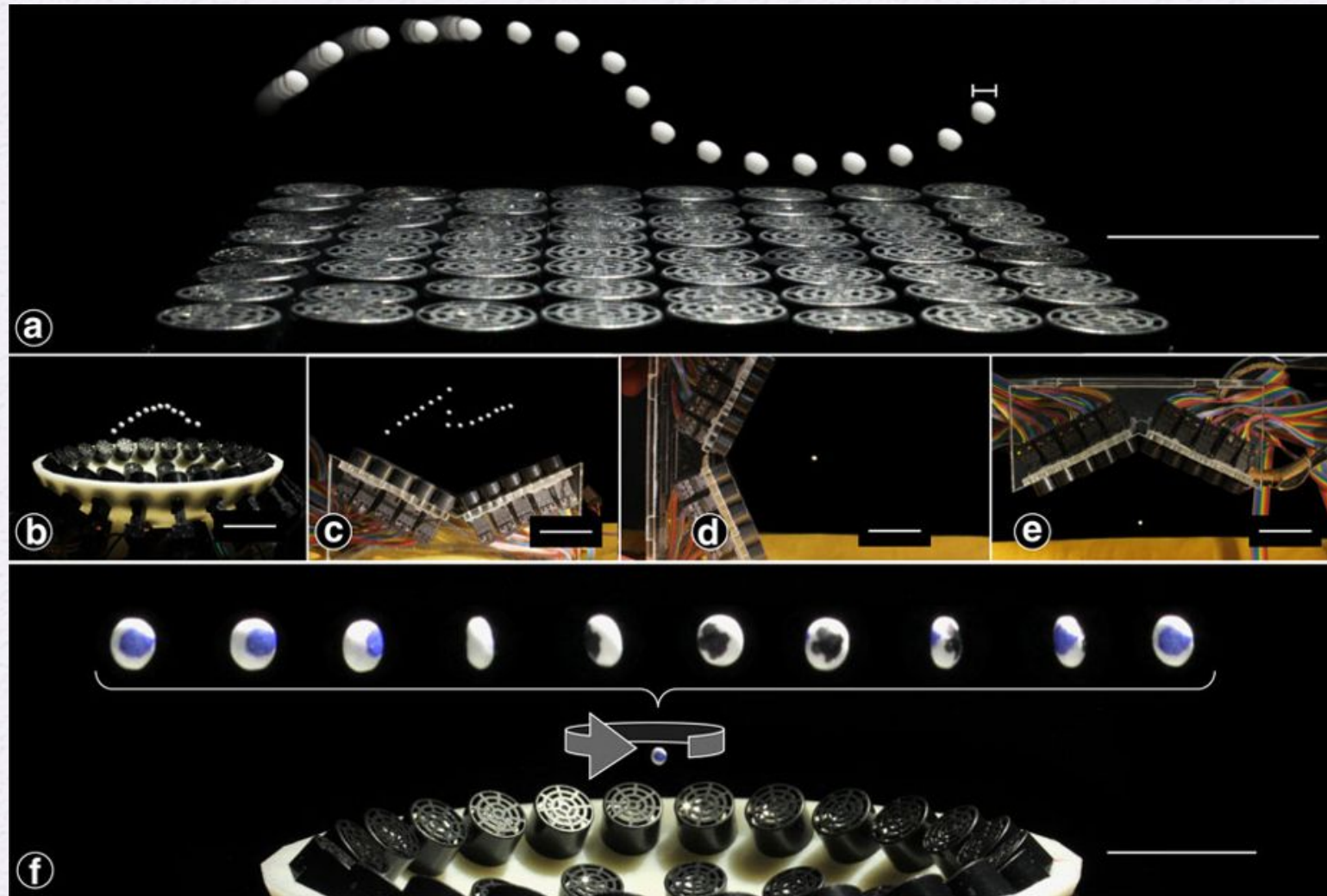
Standing wave that drifts because the left electrode emits

$$\cos\left(2\pi Ft + \int_0^t 2\pi \Delta F(\tau) d\tau\right)$$

*controled
phase shift*



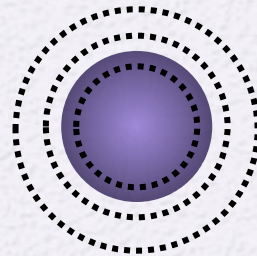
Holographic tweezers





c) On compressible bubbles: resonance and Bjernes forces

piezo

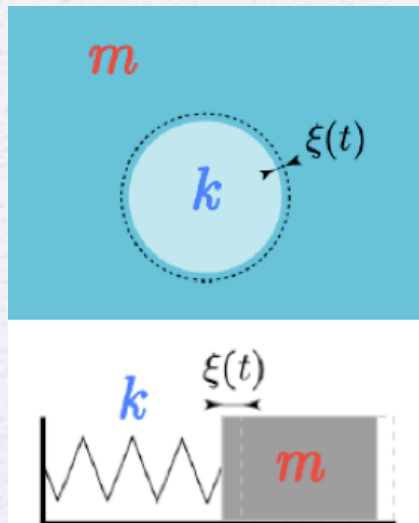


Resonance of bubbles

XVI. *On Musical Air-Bubbles and the Sounds of Running Water.* By M. MINNAERT, Sc.D., *Heliophysical Institute of the Physical Laboratory, Utrecht**.

PHYSICISTS have hardly ever investigated the sounds of running water. As a matter of fact we know very little about the murmur of the brook, the roar of the cataract, or the humming of the sea.

(*Philosophical Magazine*, 1933)



bubble = harmonic oscillator

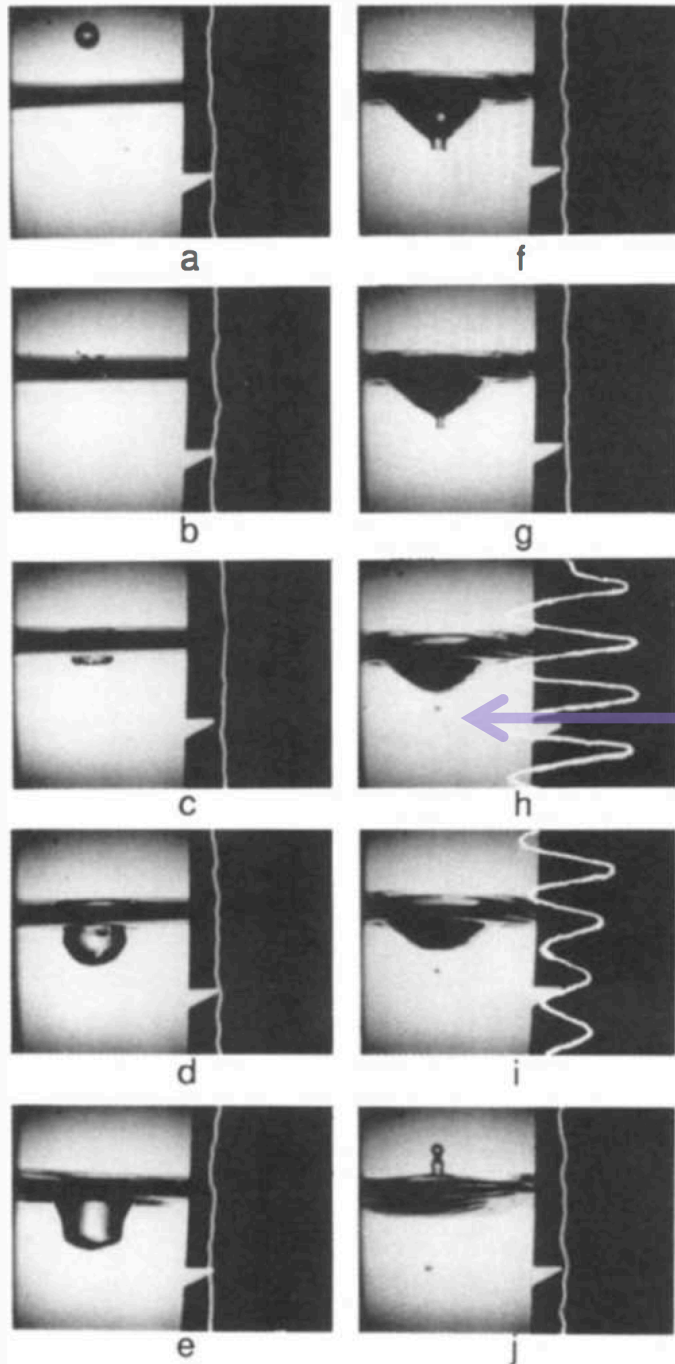
mass : displaced liquid
stiffness : gas compressibility

Minnaert formula

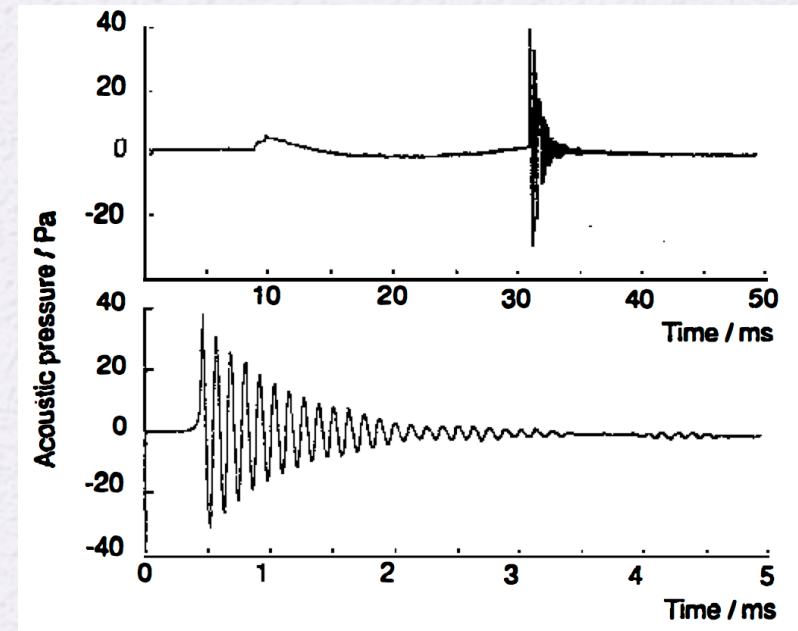
$$f \times R \simeq 3 \text{ m/s}$$

3 kHz for 1 mm bubbles

Illustrations: Olivier Vincent

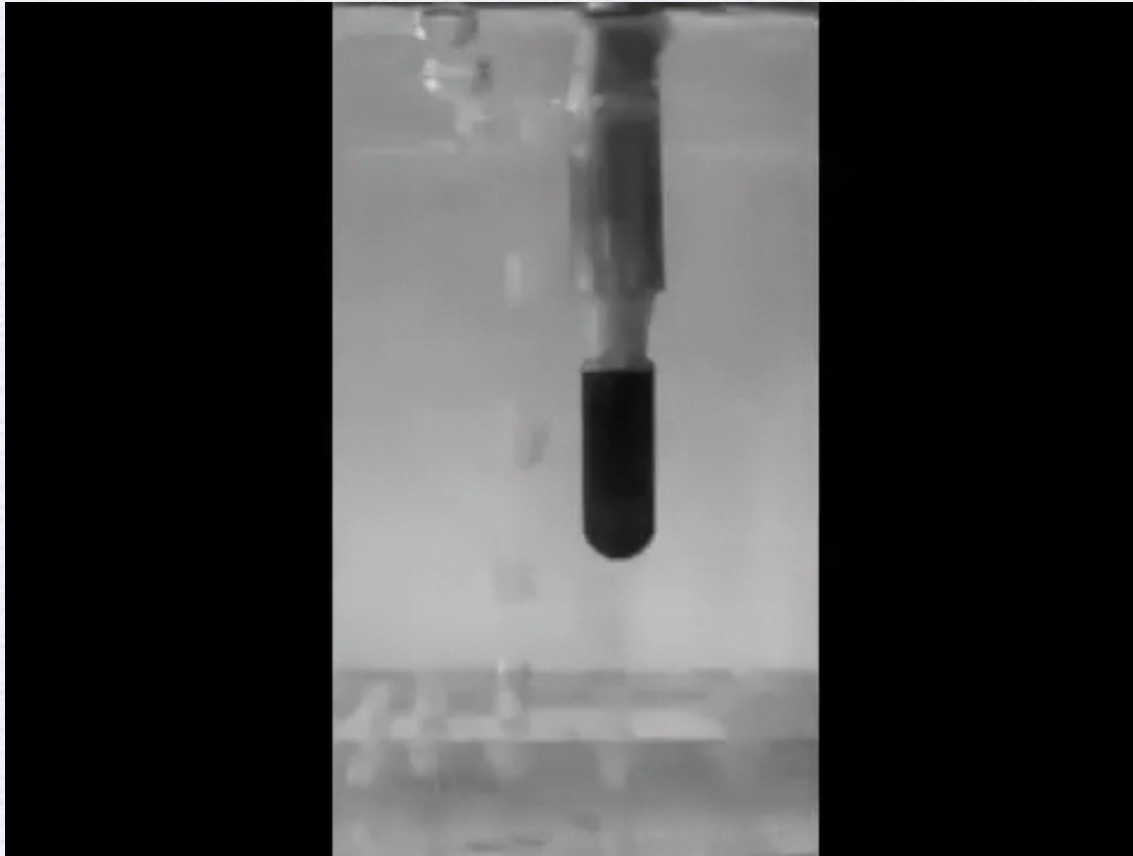


Sound of rain droplet:
due to a bubble!



← bubble

The bubble piano



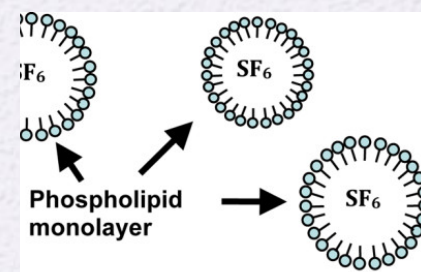
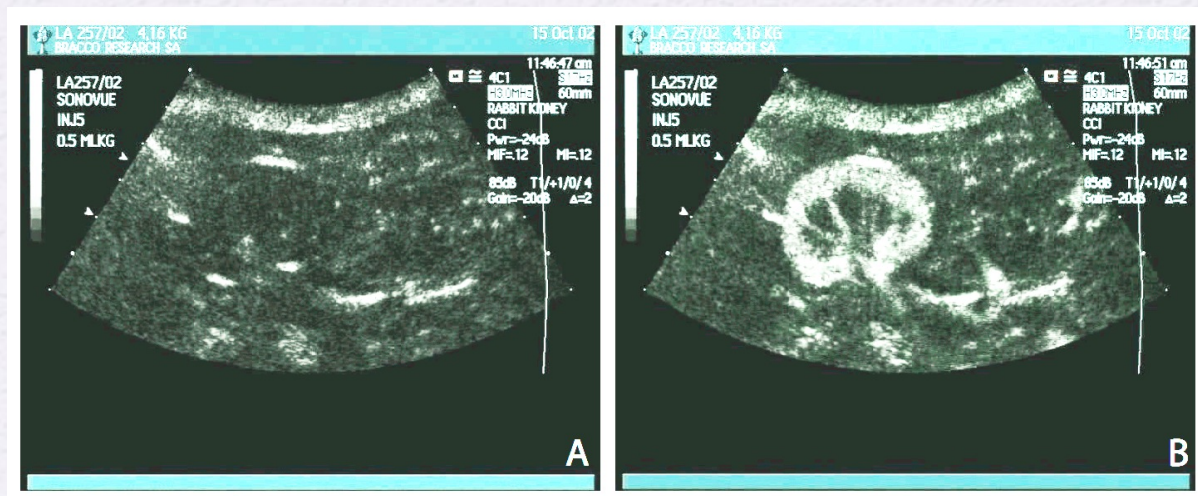
Valentin Leroy, MSC, Paris 7



Playing "Ode to Joy" with the bubble piano

Biomedical applications

- Contrast agents for ultrasound echography



Sonovue® bubbles $d \sim 3 \mu\text{m}$
Driven at $\sim 1 \text{ MHz}$

- **Sonoporation**

Near membranes enhanced permeabilisation

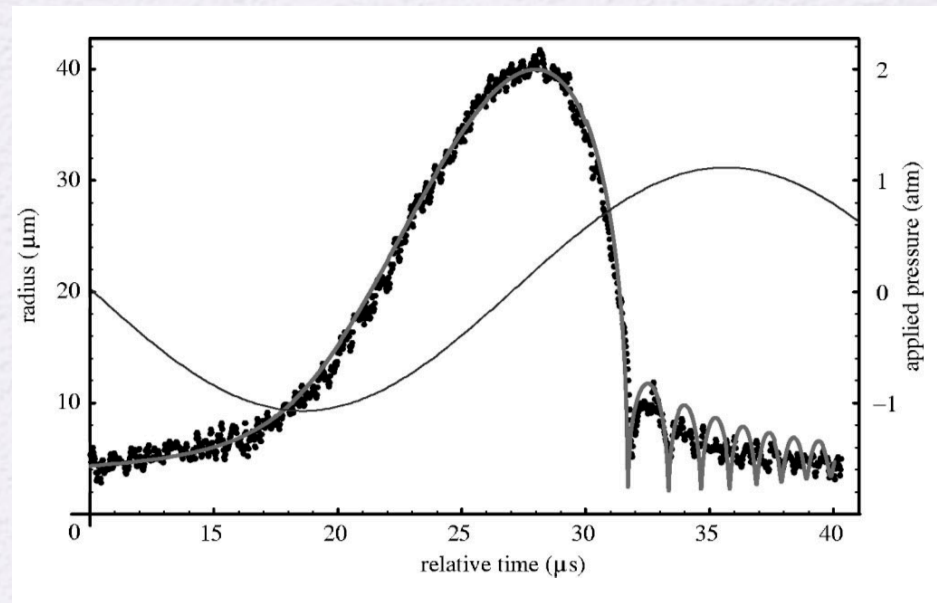
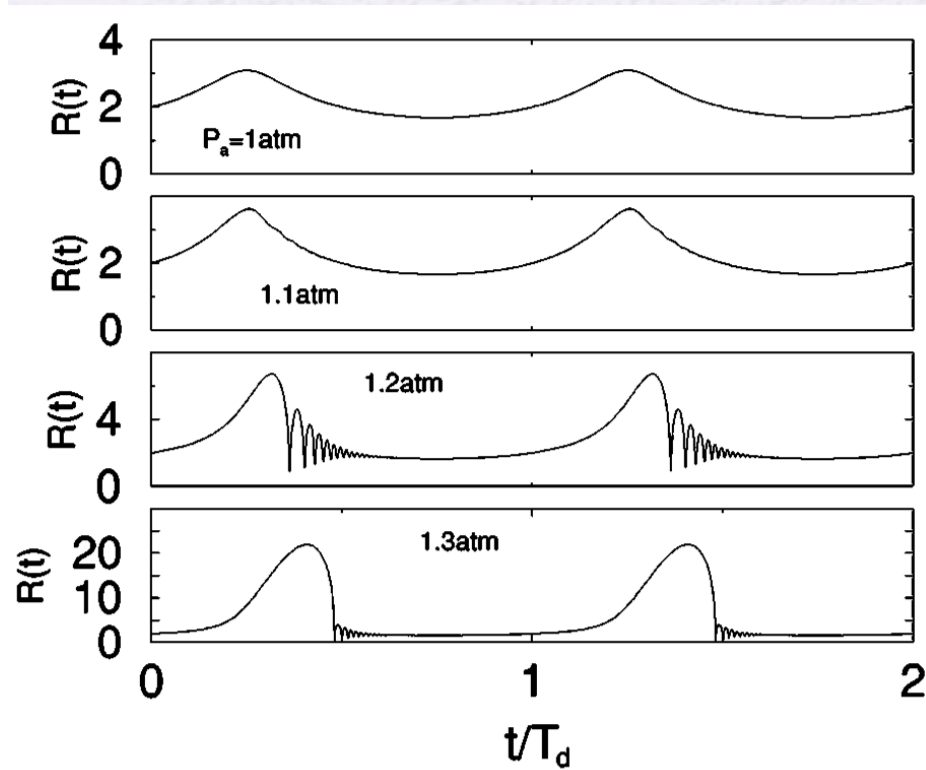
Targeted delivery



*Bracco
Research*

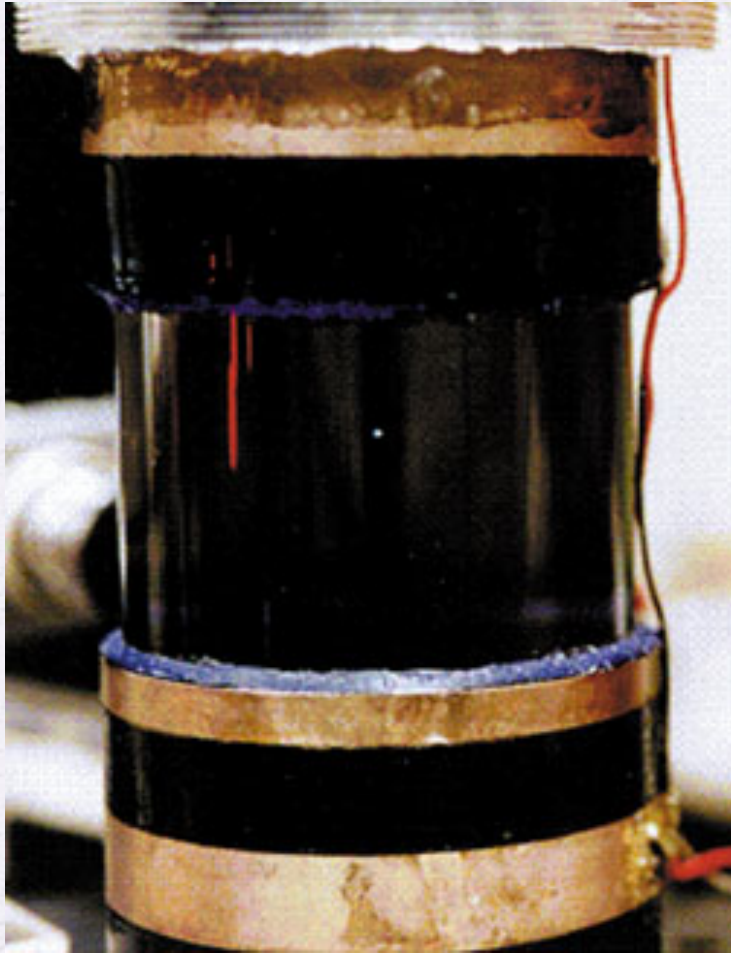
*Taniyama et al
Circulation 105
(2002)*

Excitation at large amplitude, $f \ll f_{\text{res}}$

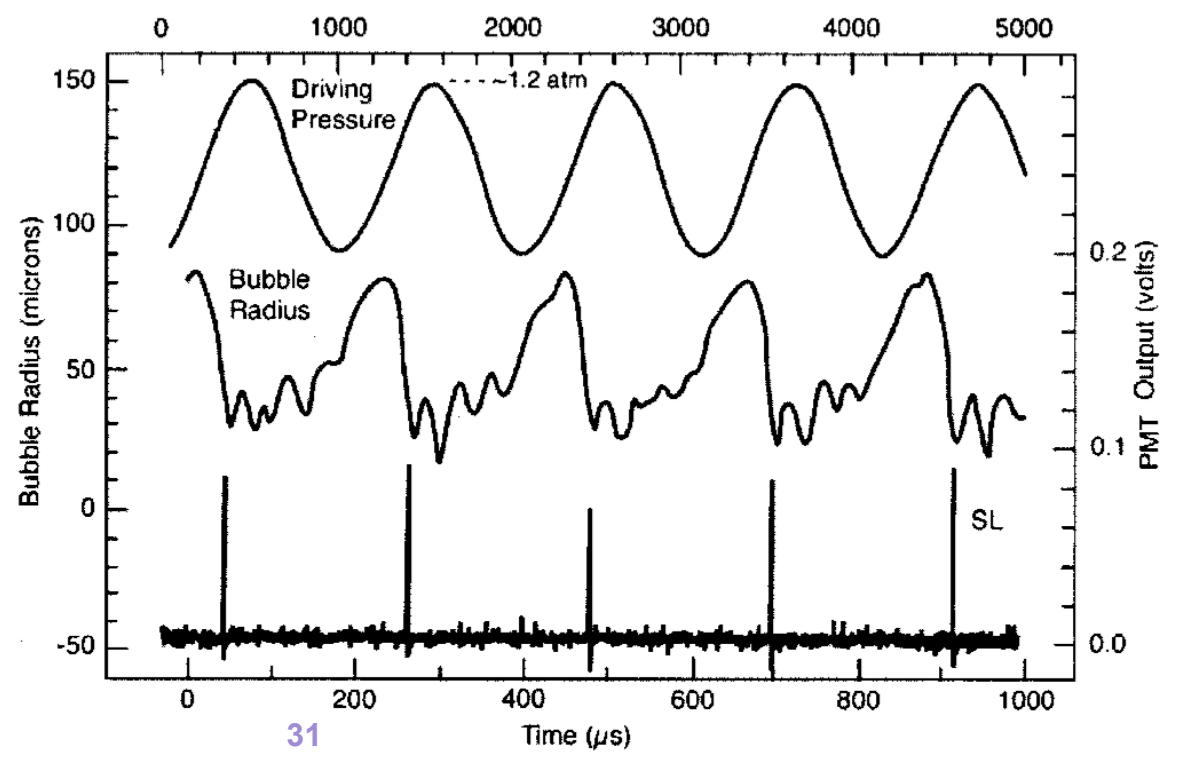
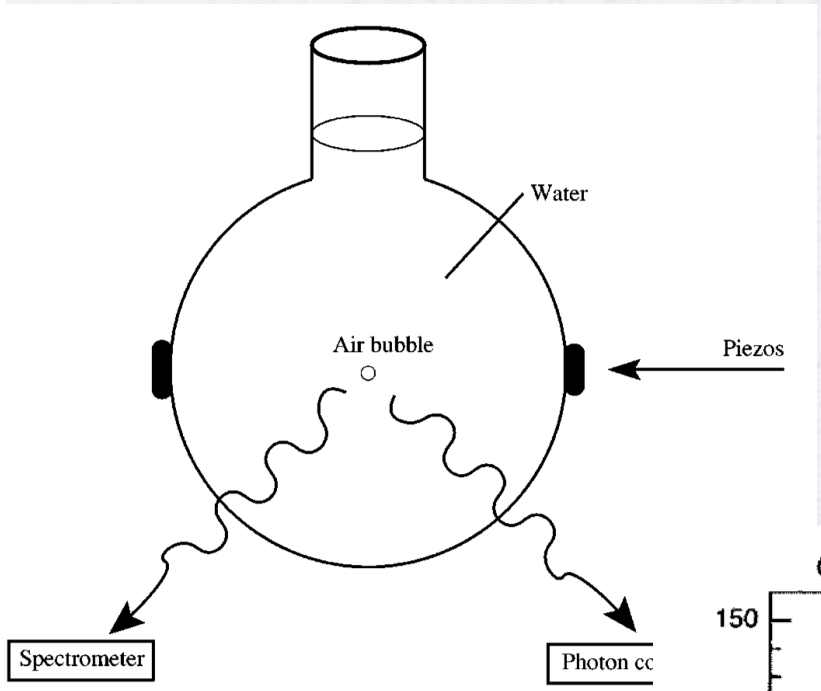


Brenner 2002

Sonoluminescence

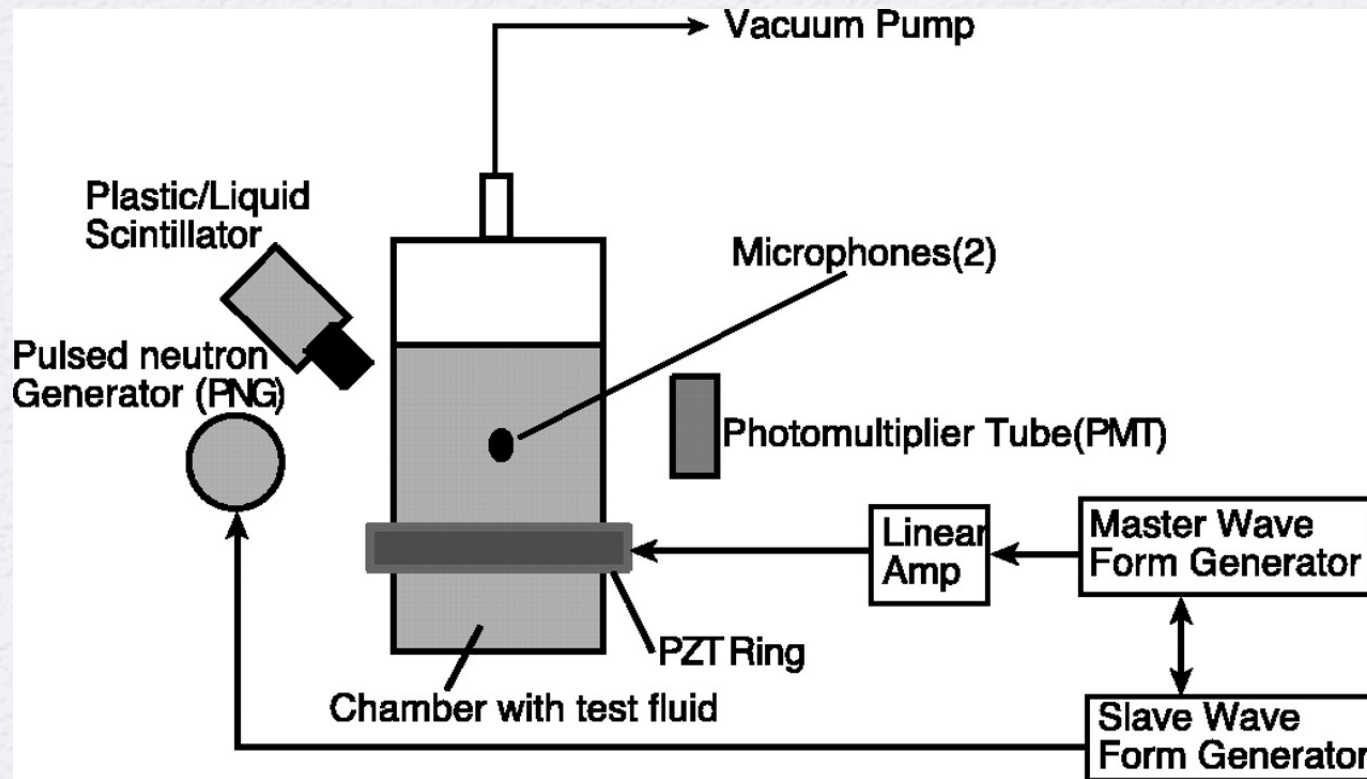


Crum 1994

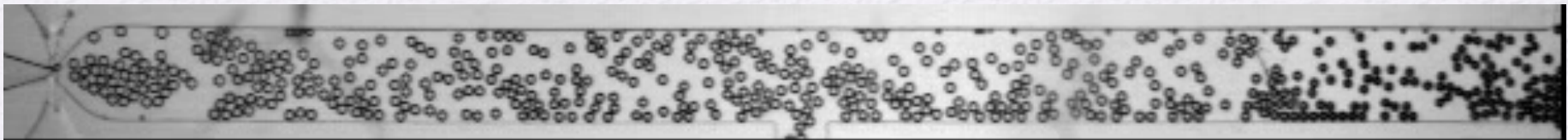
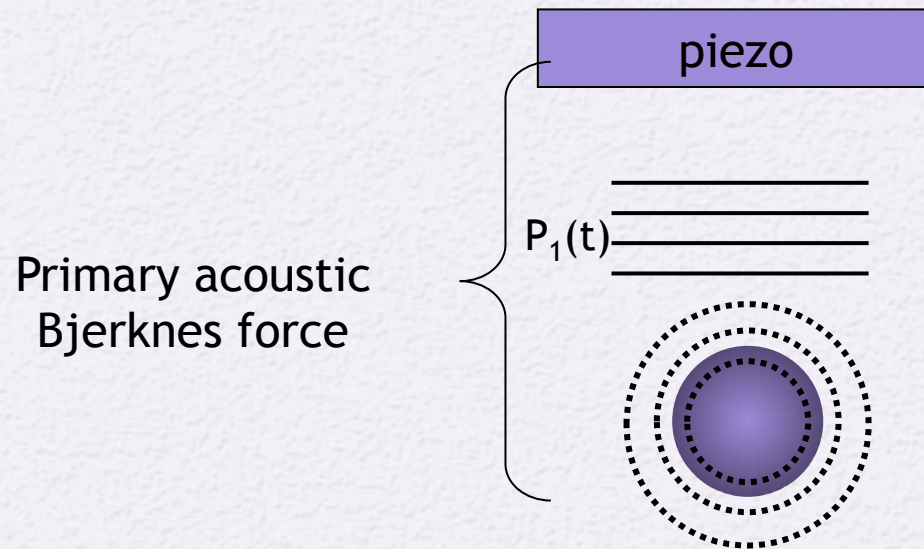


Crum 1994

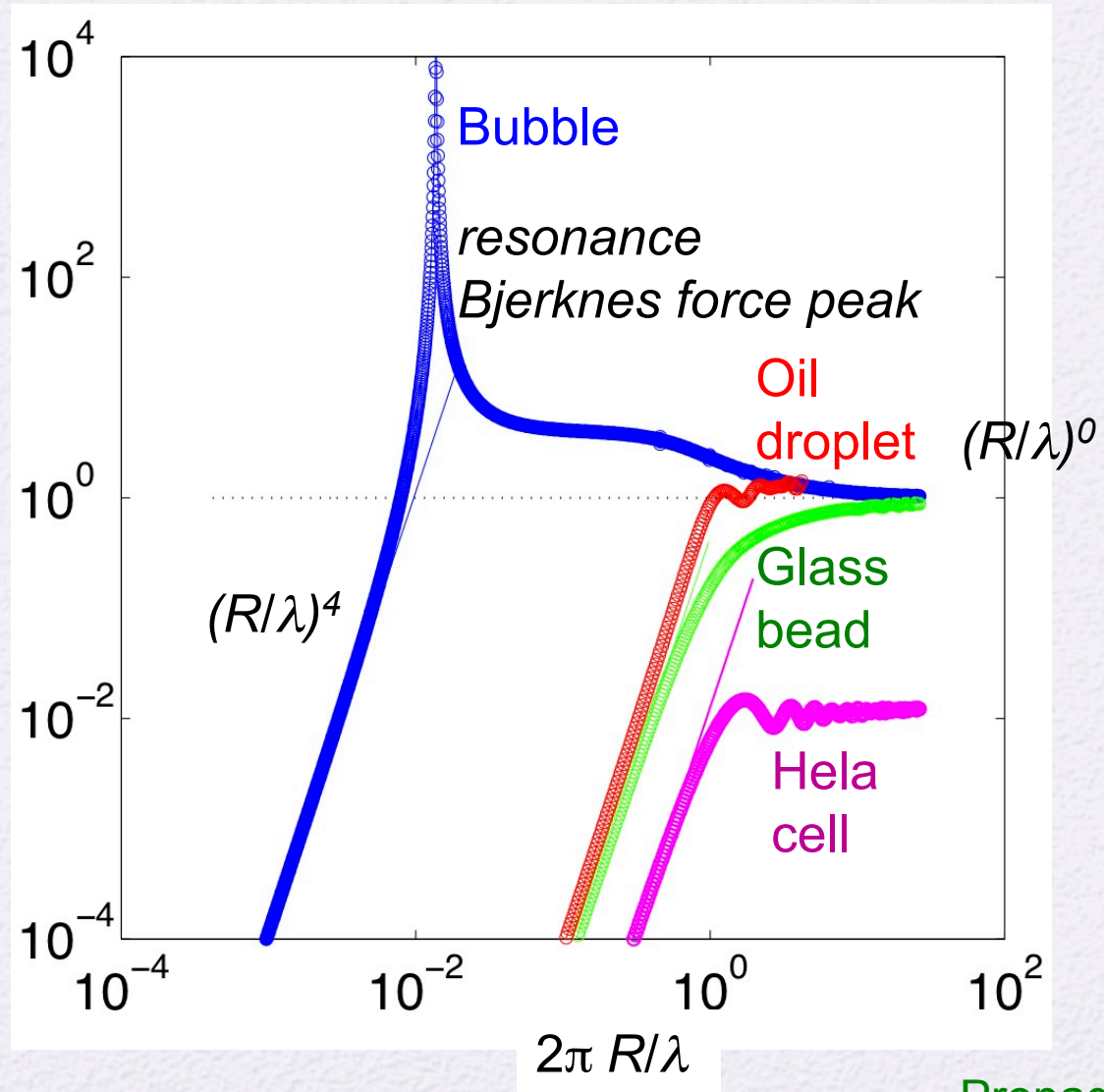
Sonofusion?



c) On compressible bubbles: resonance and Bjerknes forces



Radiation force
 F_{rad}
in units of
 $\pi R^2 I / c$



Propagating waves
Yosioka 1955

The Bjerknes force

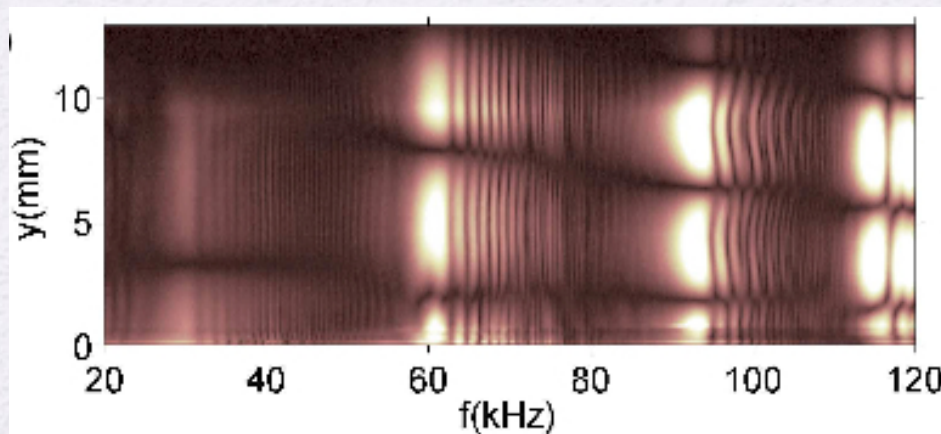
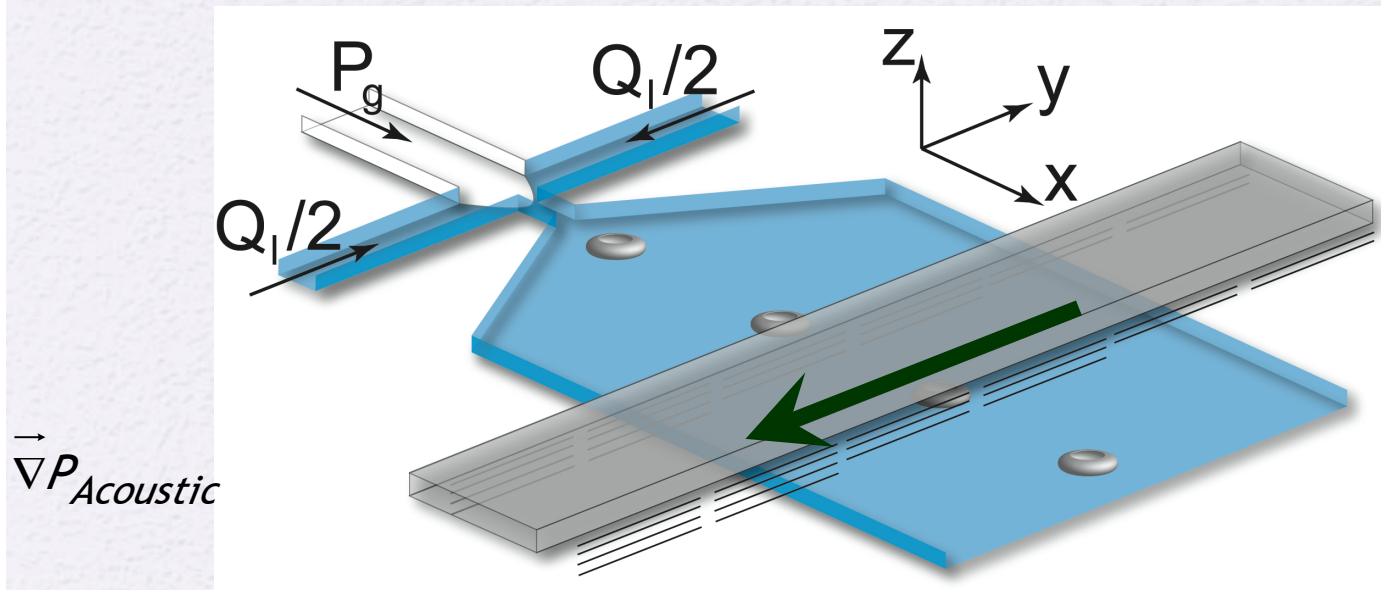
- Bjerknes force:
due to large **volume oscillation** at resonance,
in a pressure gradient

$$\begin{aligned} F_{\text{Bjerknes}} &= \left\langle -V(t) \nabla P_{\text{Acoustic}}(t) \right\rangle \\ &= \left\langle - \left[V_0 + \delta V \cos(\omega t + \varphi) \right] \left[\nabla P_A \cos(\omega t) \right] \right\rangle \\ &= -\delta V \nabla P_A \cos \varphi \end{aligned}$$



Bjerknes 1906

Set-up to induce Bjerknes forces



deviations

With a change of the standing wave pattern

$\vec{\nabla}P_{Acoustic}$



$f = 147 \text{ kHz}$

$\vec{\nabla}P_{Acoustic}$

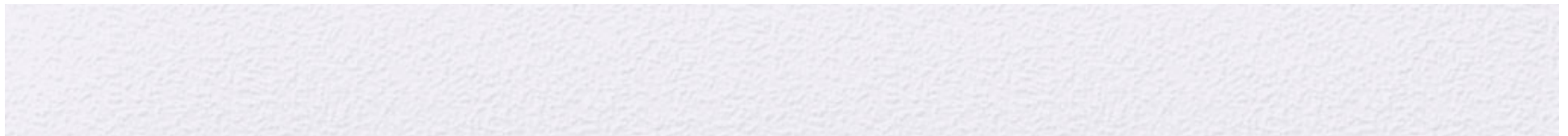
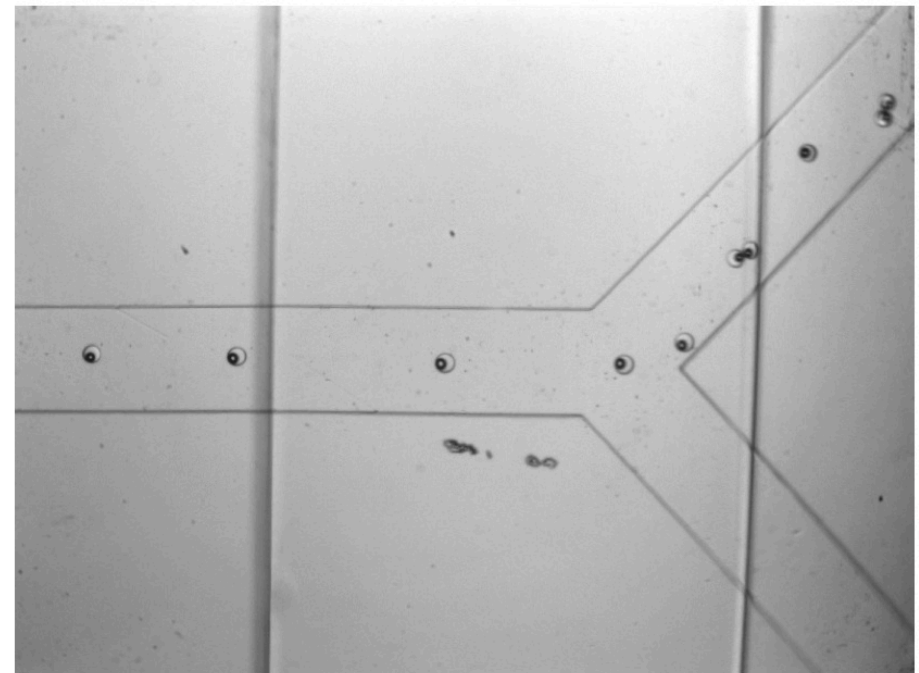
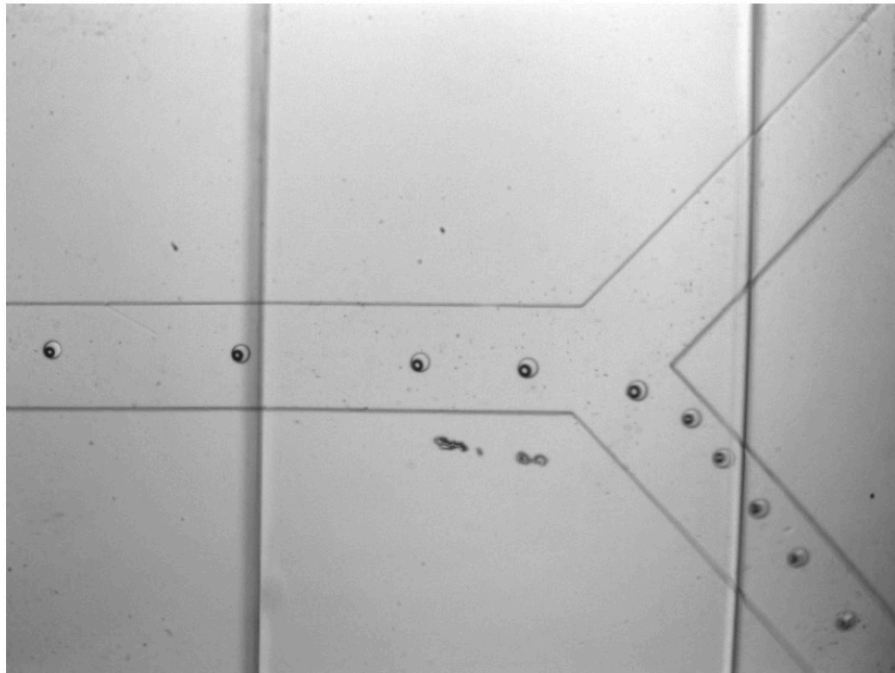
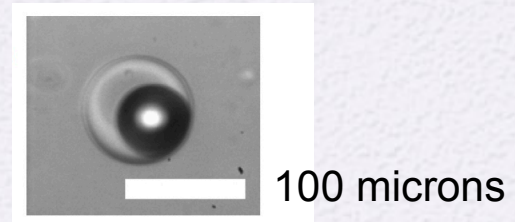


$f = 151 \text{ kHz}$

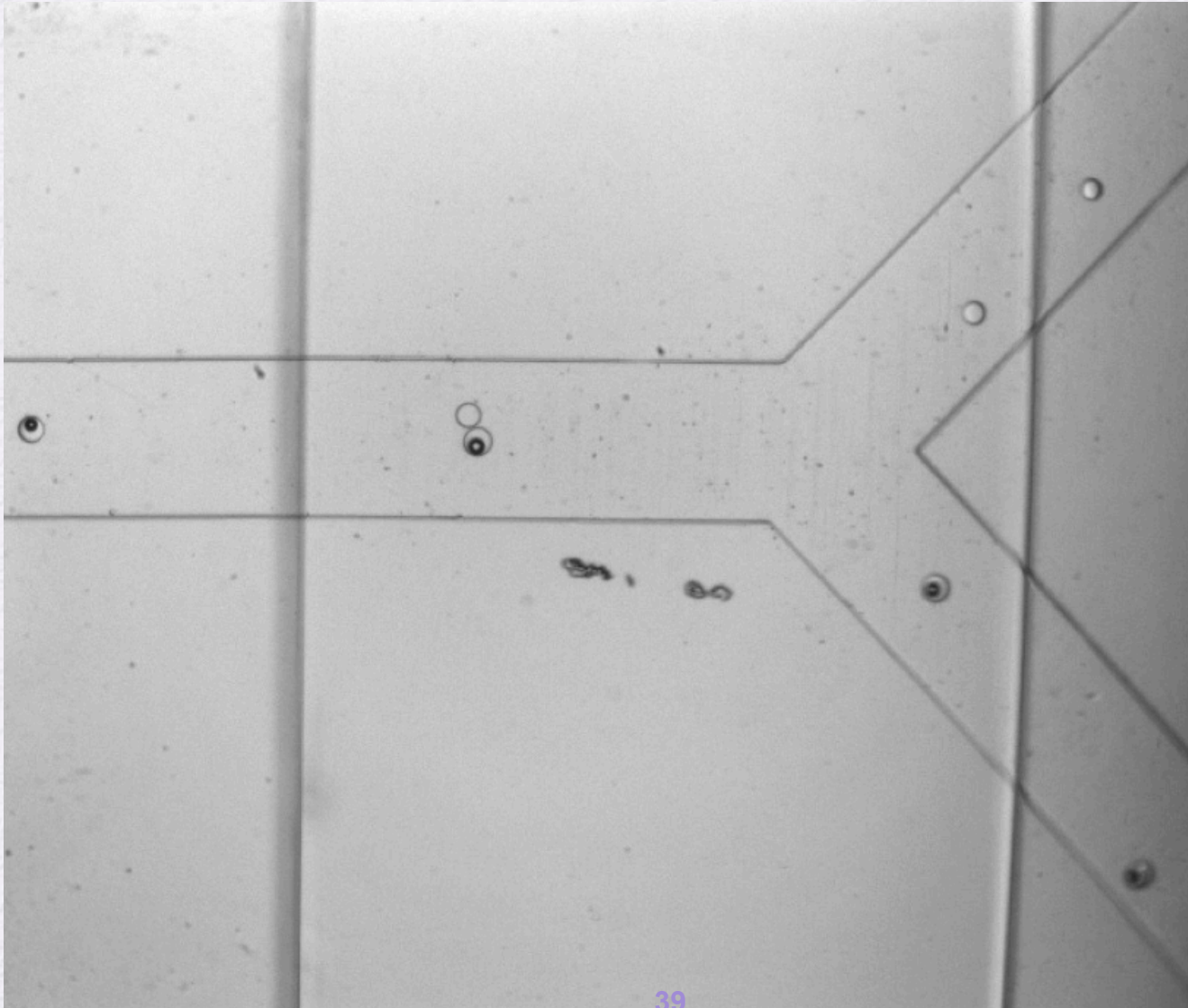
Rabaud, Thibault, Raven, Hugon,
Lacot and Marmottant
Phys. Fluids 2011

Example of applications of Bjerknes forces

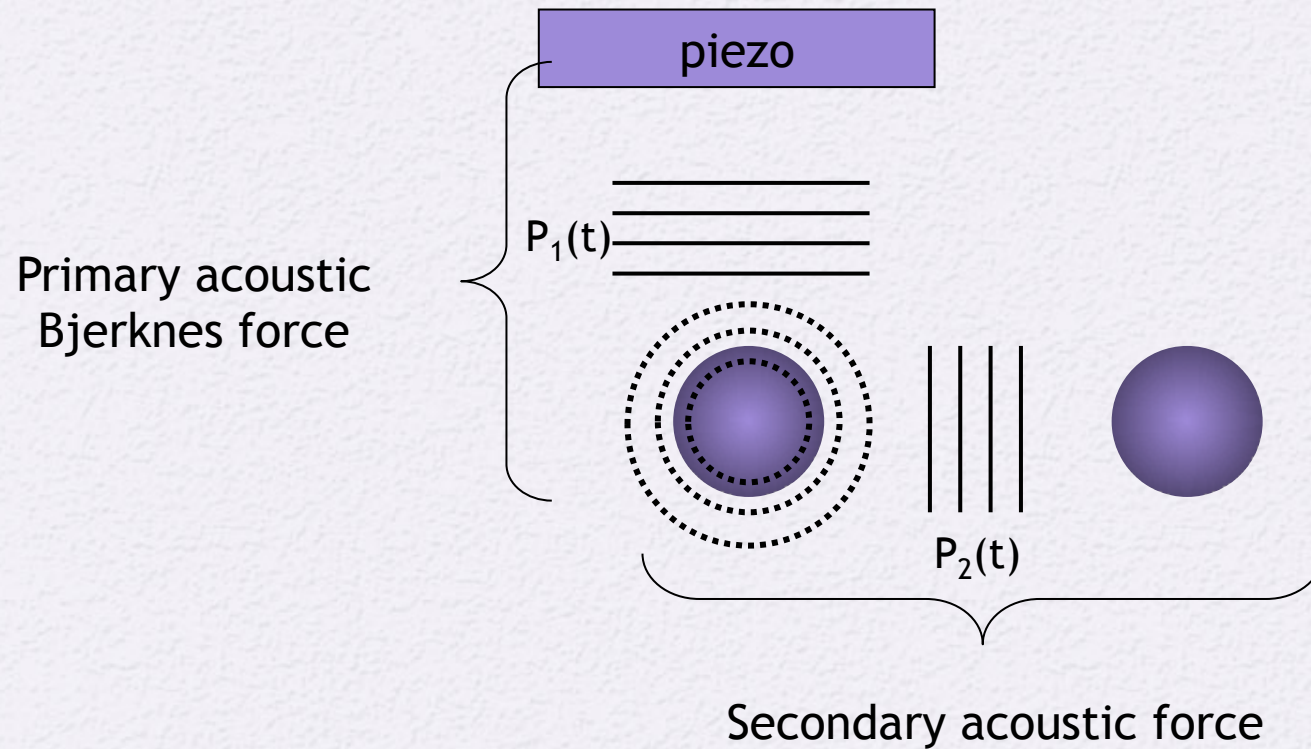
- Switch droplets containing a bubble



- Sort droplets



Many bubbles: Self-organisation



Primary acoustic
Bjerknes force

Secondary acoustic force

attractive at small amplitude *Crum 1975*

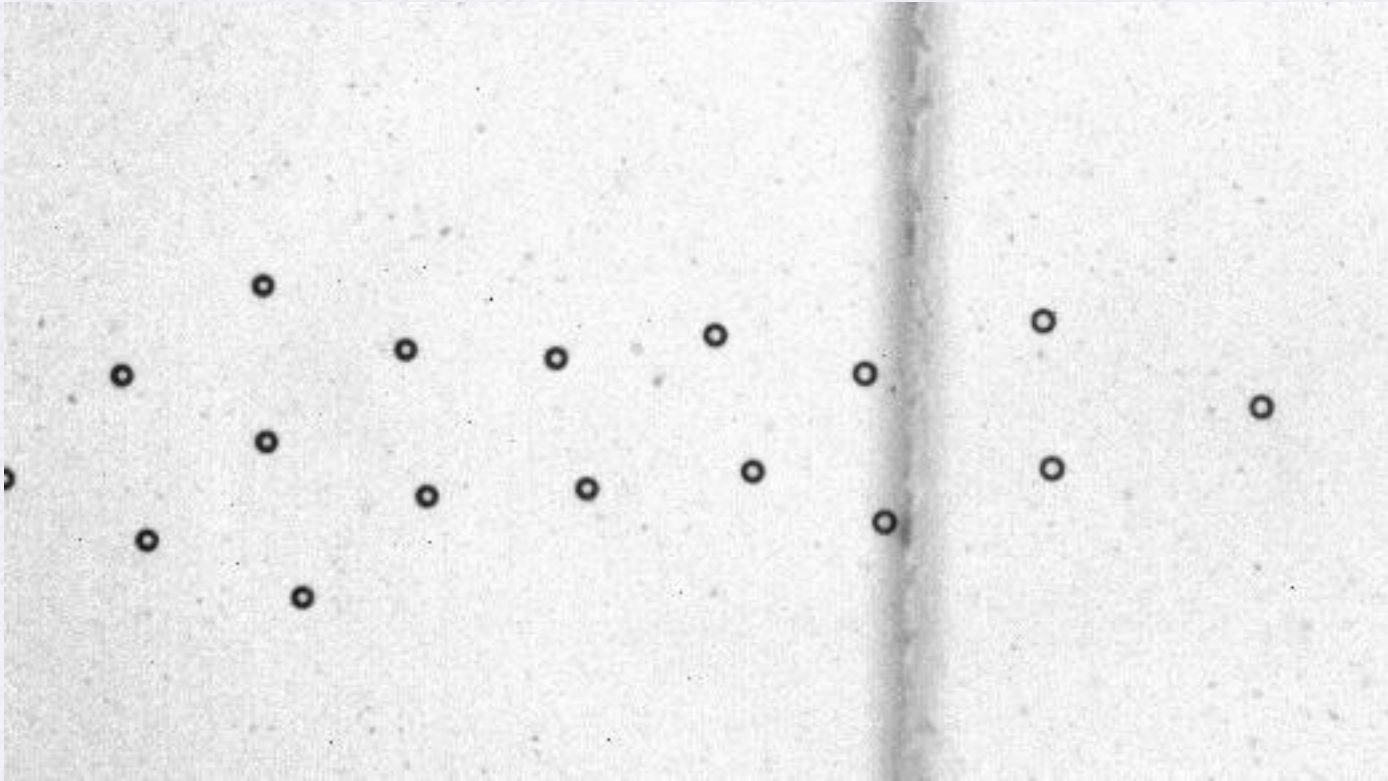


Ultrasonic "ballet" at larger amplitudes.



Bubbles keep their distance...

acoustically-bound crystals

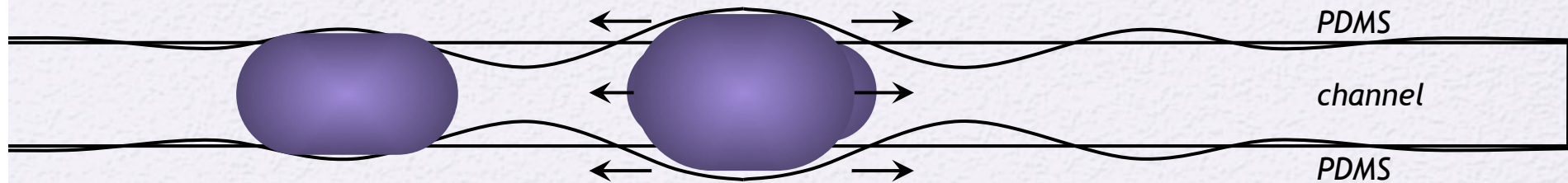


...and acoustically-bound crystals are formed.

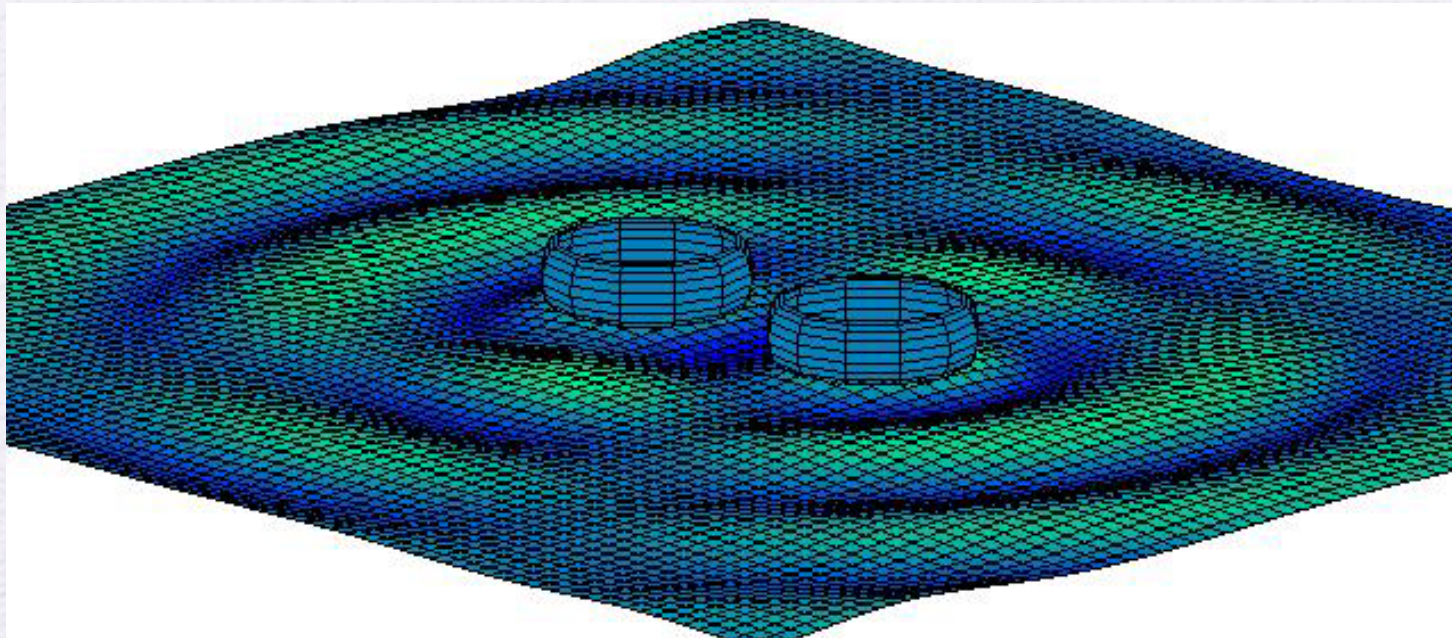
Rabaud, Thibault, Mathieu and
Marmottant *Phys. Rev. Lett.* 2011

Interaction through surface acoustic waves, on PDMS!

Side view of a pulsating bubble

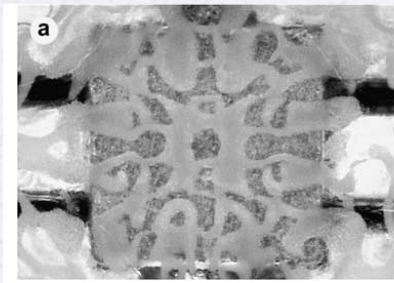
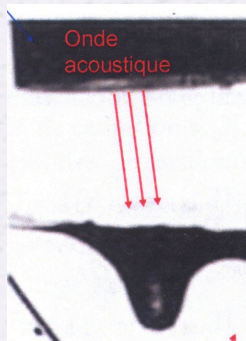


Surface acoustic waves (Rayleigh waves) on the PDMS walls



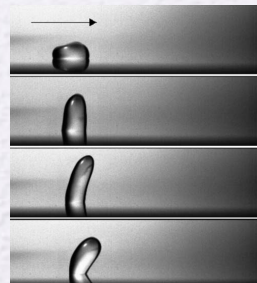
Outline

1) Radiation pressure



- a) on interfaces
- b) on particles, tweezers
- c) on bubbles

2) Acoustic streaming



- a) Quartz wind
- b) near boundaries
- c) near bubbles

2) Acoustic streaming

Navier-Stokes equation

$$\rho \frac{\partial \mathbf{v}}{\partial t} = -\nabla p + \eta \nabla^2 \mathbf{v} - \rho (\mathbf{v} \cdot \nabla) \mathbf{v}$$

Inertia

Consequences of inertia

- 1) non-reversibility of fluid (ex: flamme)
- 2) turbulence
- 3) non-linear, steady forcing on average:

Navier-Stokes averaged over time

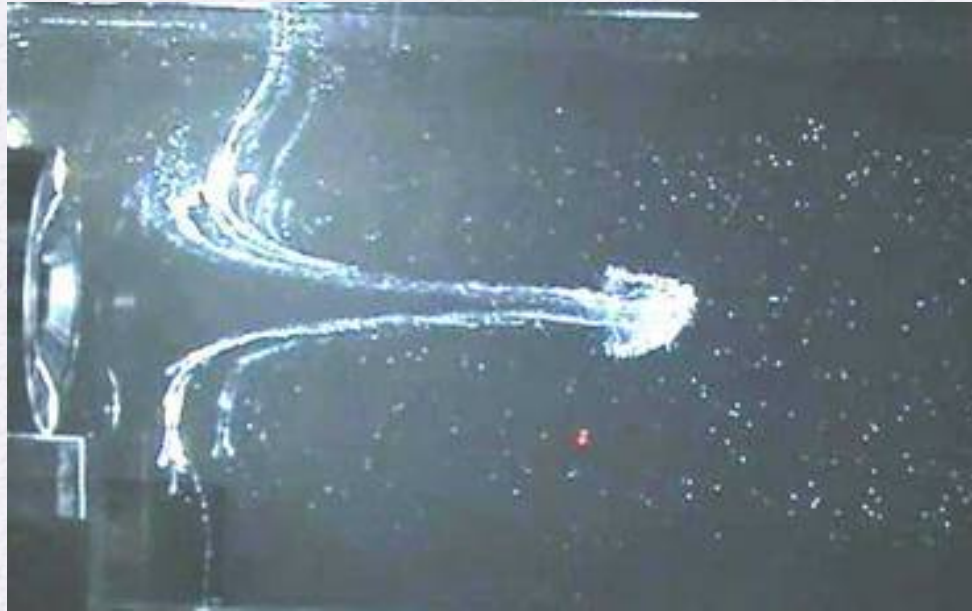
$$\mathbf{0} = -\nabla \langle p \rangle + \eta \nabla^2 \langle \mathbf{v} \rangle - \langle \rho (\mathbf{v} \cdot \nabla) \mathbf{v} \rangle$$

steady streaming

Reynolds stress

$$\mathbf{f}_{\text{Reynolds}} = \langle -\rho (\mathbf{v} \cdot \nabla) \mathbf{v} \rangle \quad \text{quadratic}$$

a) Quartz wind



Attenuated wave

$$\mathbf{v} = \mathbf{v}_A \exp(-\alpha x) \cos(\omega t - kx)$$

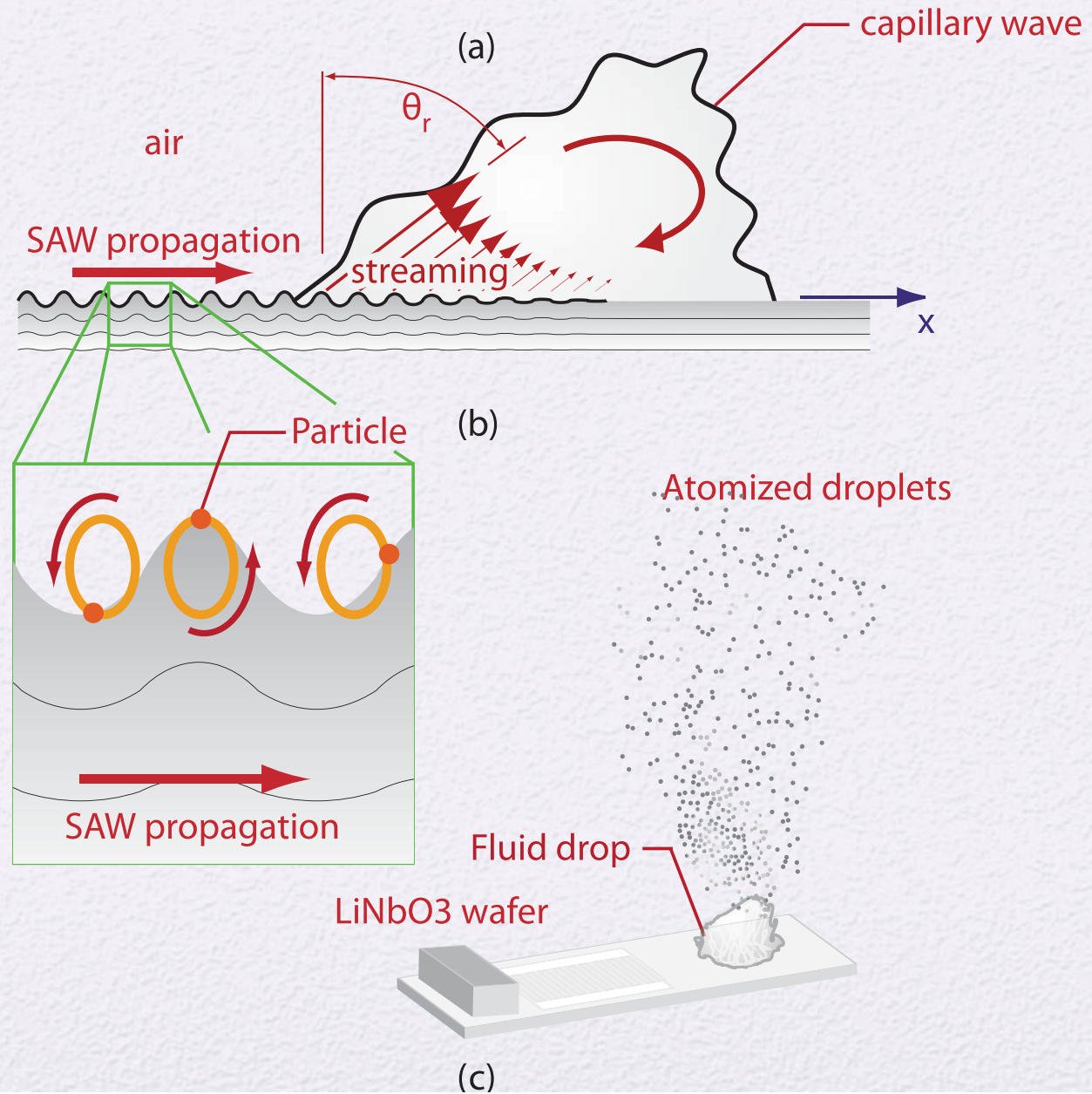
$$\mathbf{f}_{\text{Reynolds}} = \frac{1}{2} \rho \alpha v_A^2 \exp(-2\alpha x) \mathbf{e}_x$$

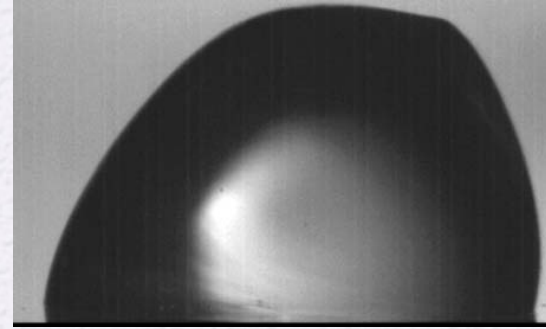
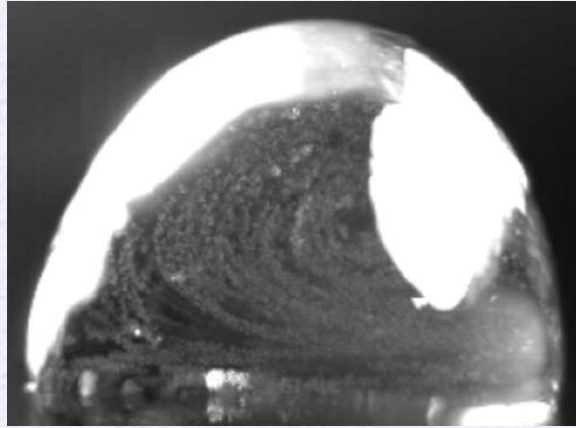
Attenuation length in water $\alpha^{-1} = 50 \text{ m.MHz}^2 / f^2$

In air $\alpha^{-1} = 1 \text{ m.MHz}^2 / f^2$

Botton

In a droplet

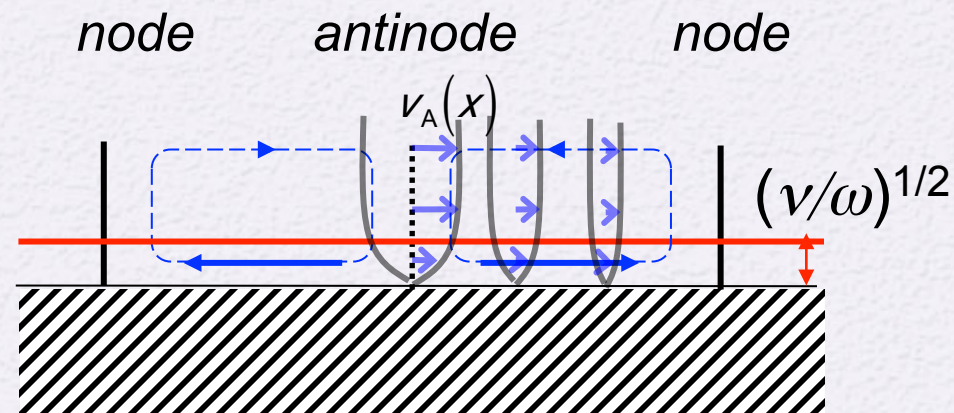
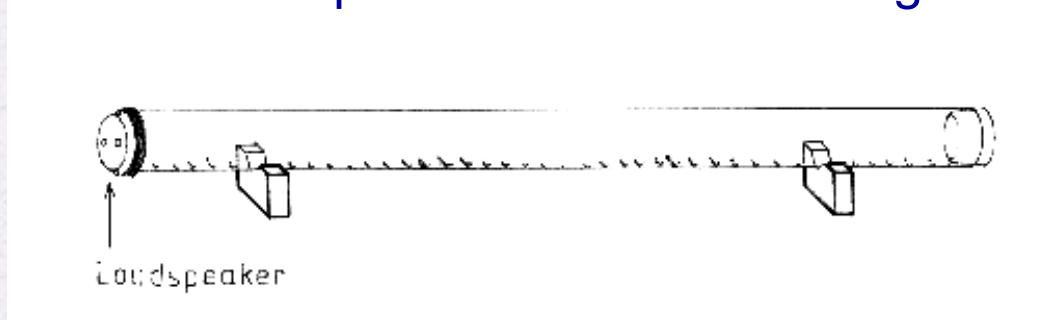




Brunet 2010

b) near a wall

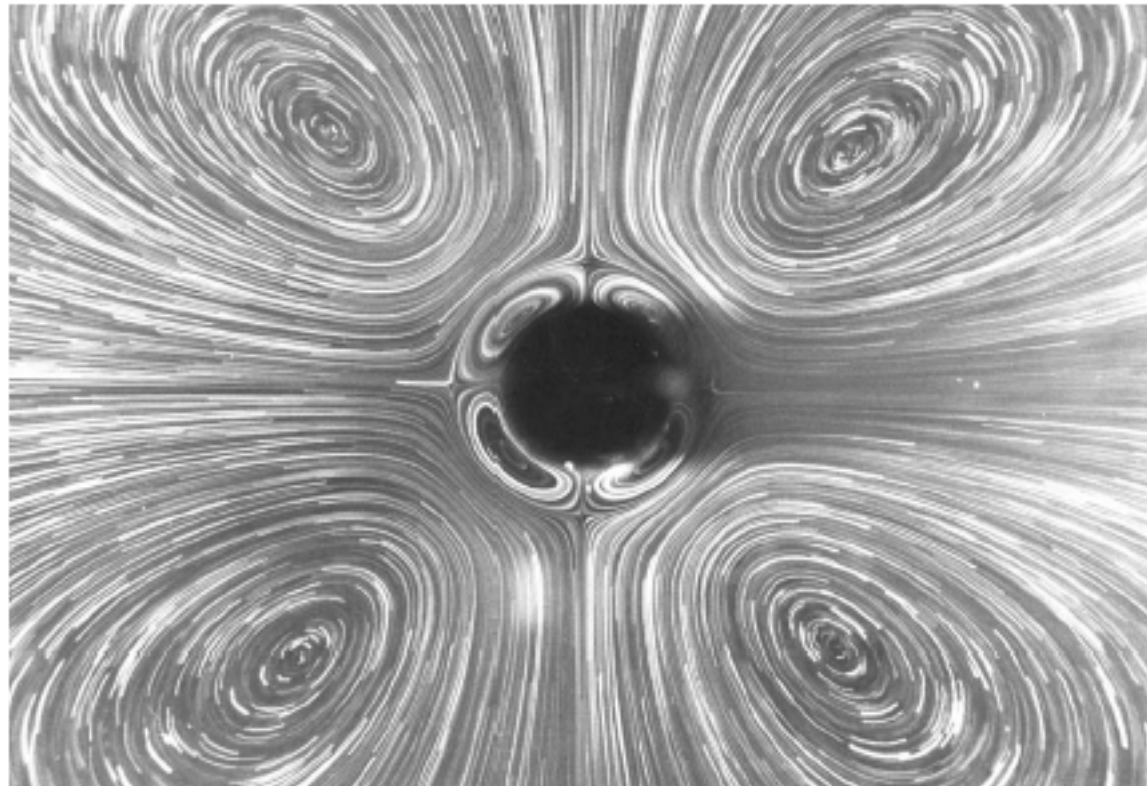
A very well known example of acoustic streaming: the Kundt tube (1866)



Reynolds stresses are important in the Stokes oscillatory boundary layer

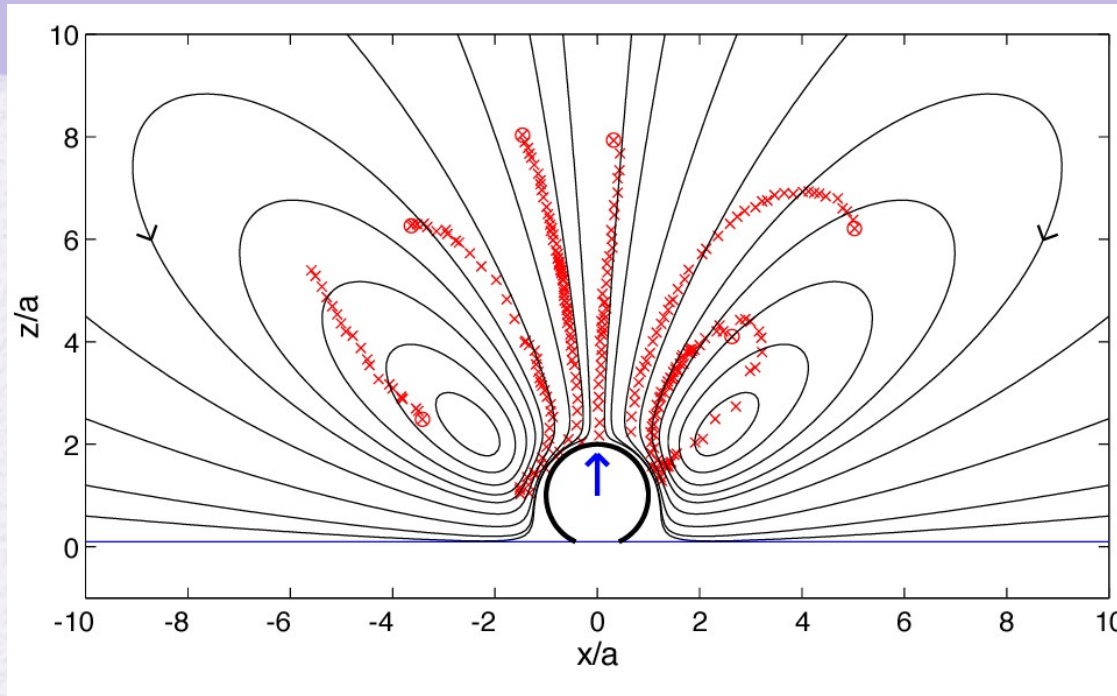
$$\mathbf{f}_{\text{Reynolds}} = \langle -\rho(\nabla \cdot \mathbf{v})\mathbf{v} \rangle \sim -\rho V_A \frac{\partial V_A}{\partial x} \mathbf{e}_x$$

b) near a wall: vibrating cylinder



Tatsuno 1980

c) Near bubbles

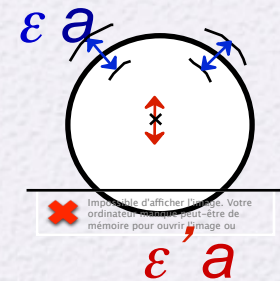


Gentiles oscillations de la bulle $\varepsilon \ll 1$

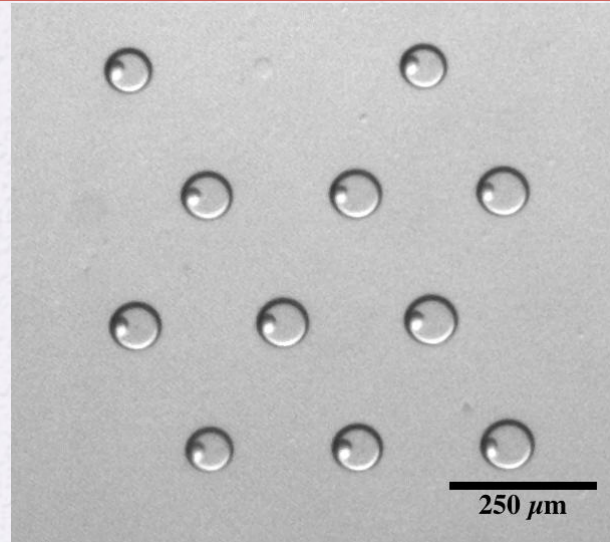
- **rayon** $r = \varepsilon a \cos(\omega t)$
- **centre** $z = \varepsilon' a \cos(\omega t - \phi)$

Streaming

$$\langle u_2 \rangle \sim \varepsilon \varepsilon' a \omega \sin \phi$$

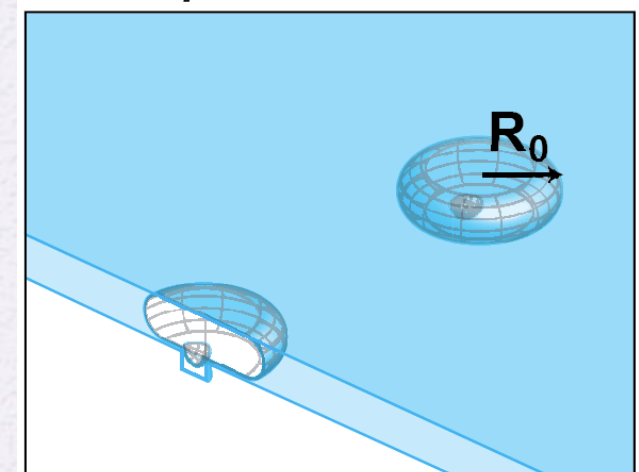


Microfluidic set-up: flattened and anchored bubbles

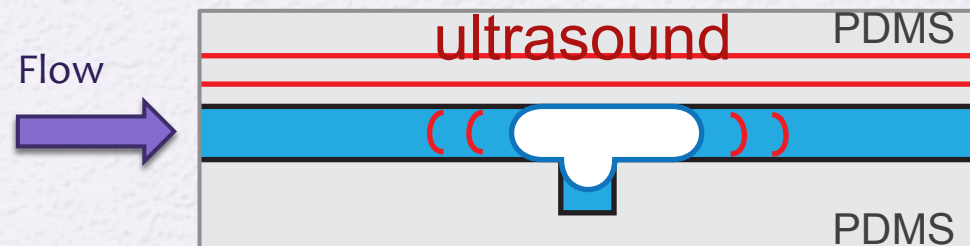


Top view

Perspective



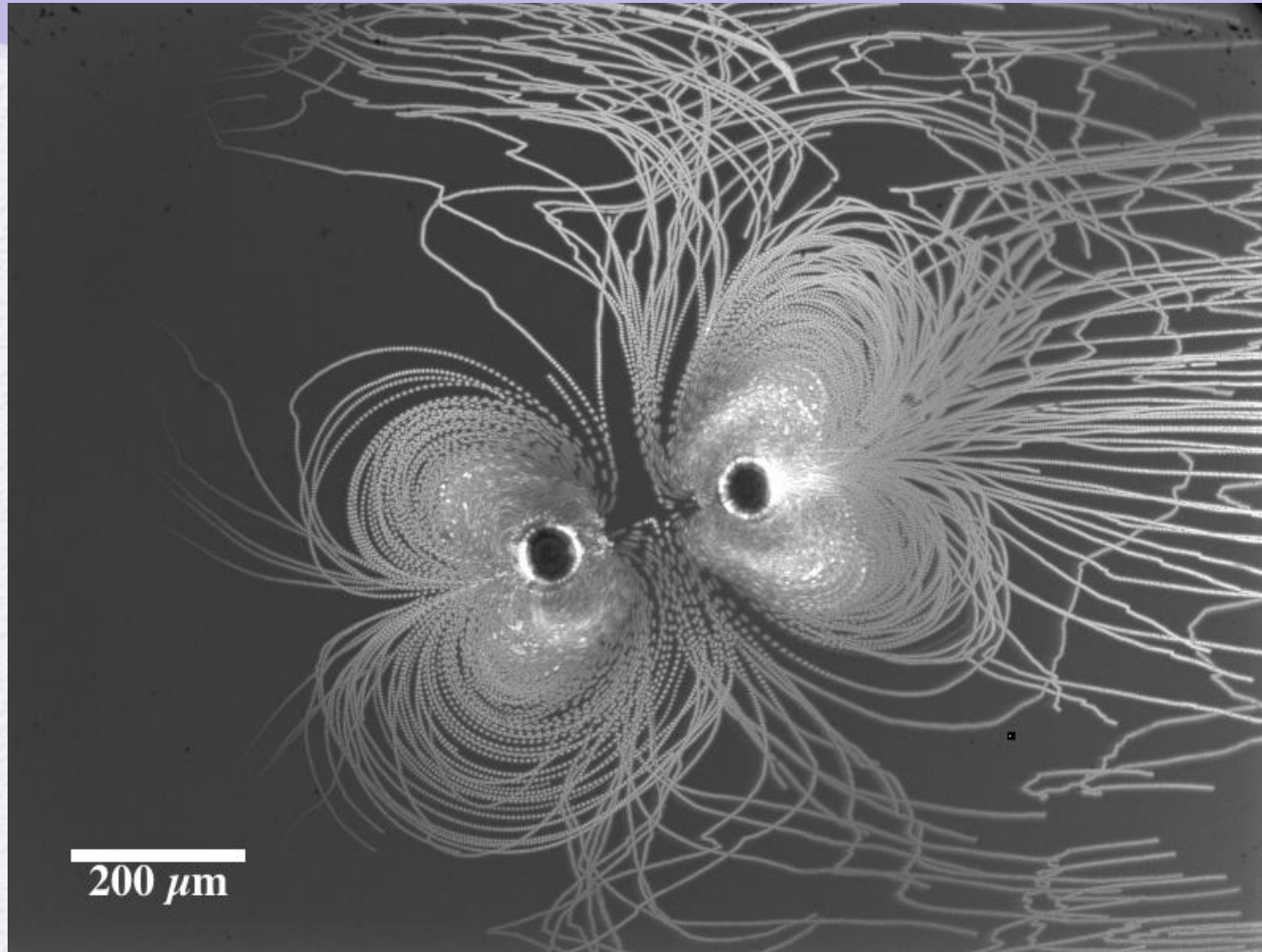
Cross-section



anchors: micropits

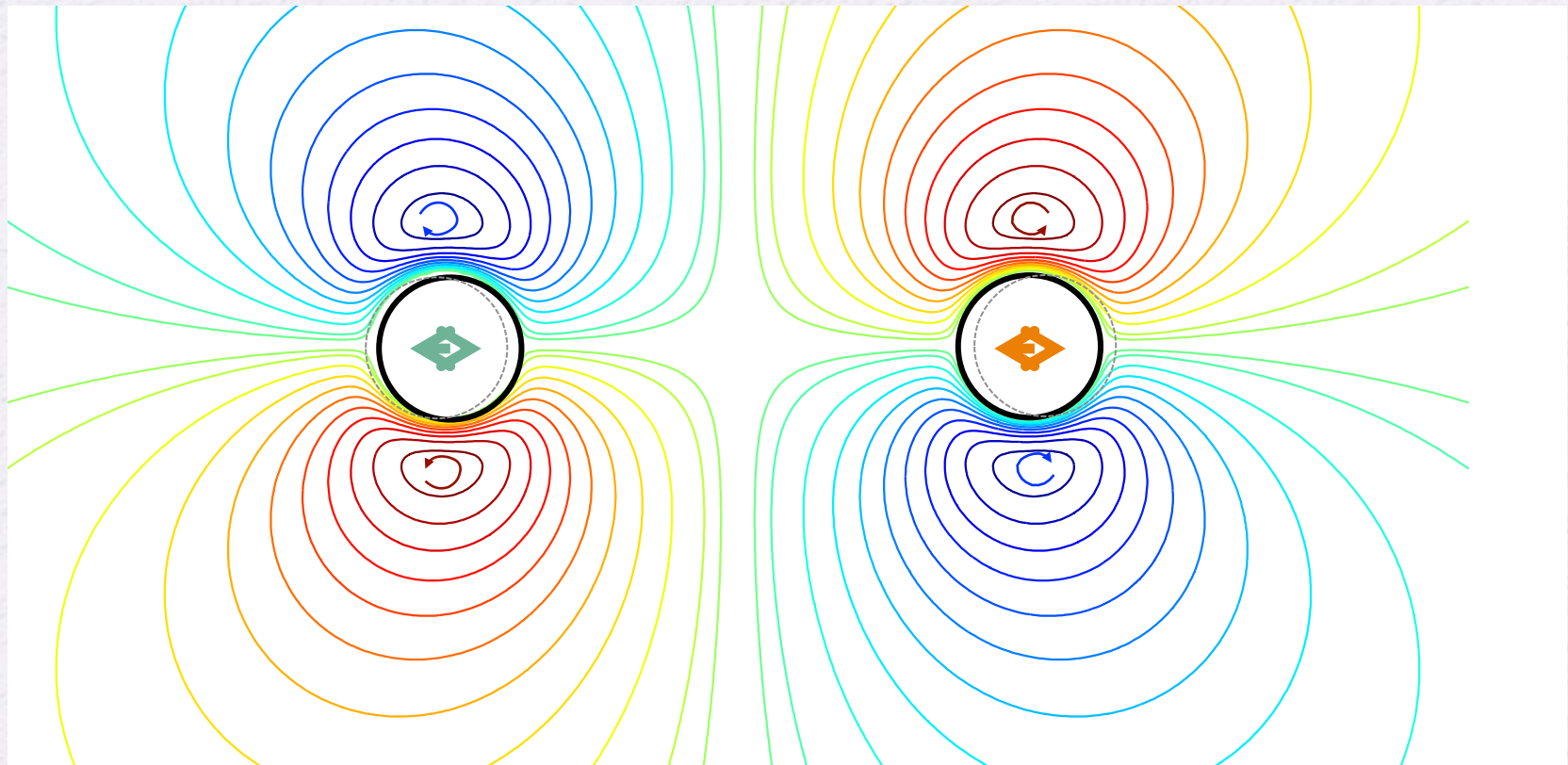
Pit method by Dangla et al 2011

Two bubbles vibrating: passionate streaming!



Acoustic streaming:

predictions from translation+oscillation vibrations



$$\Psi_L = \Psi_L^{\text{bubble 1}} + \Psi_L^{\text{bubble 2}}$$

- Seven bubbles



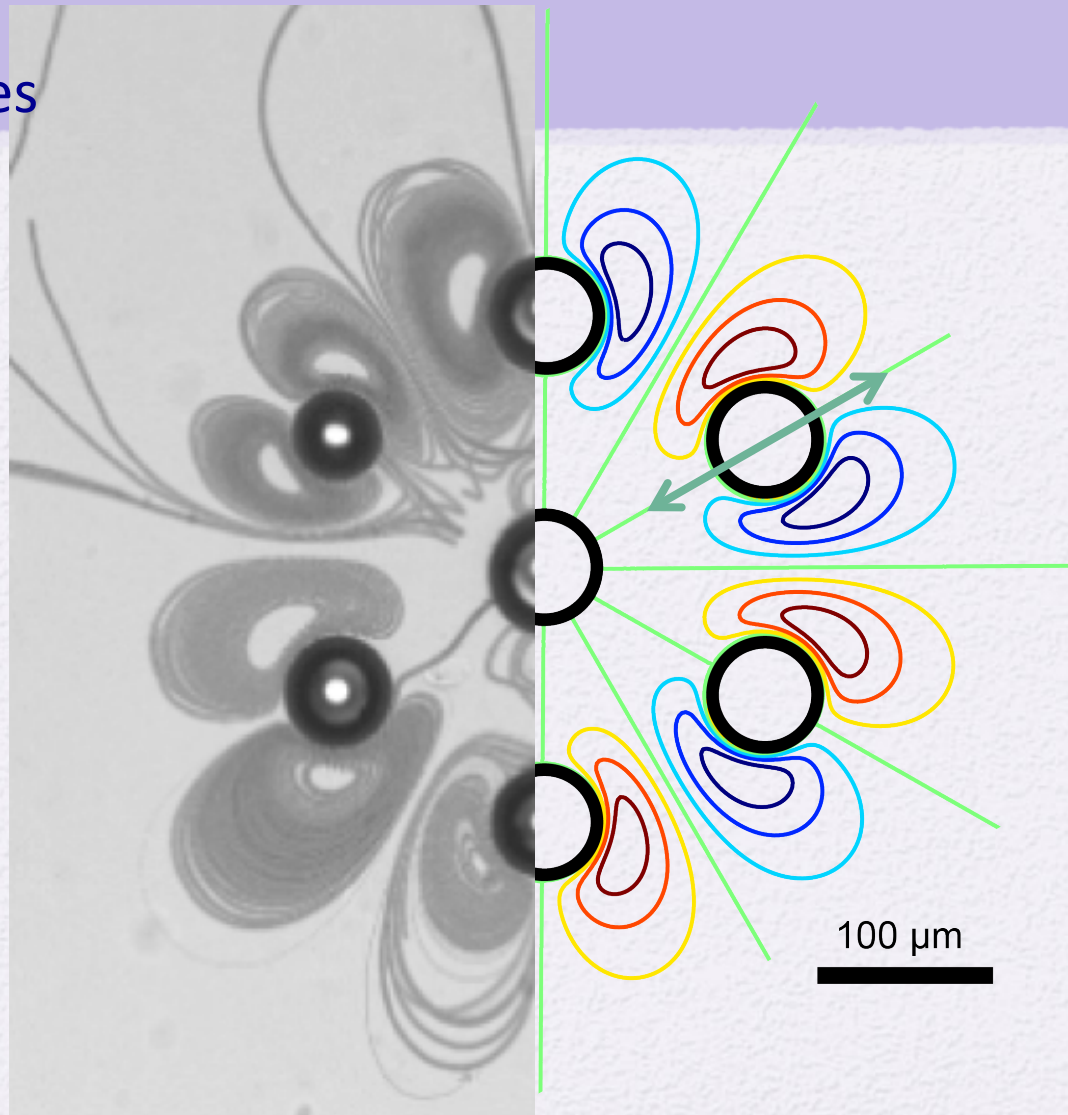
Neighbour 1

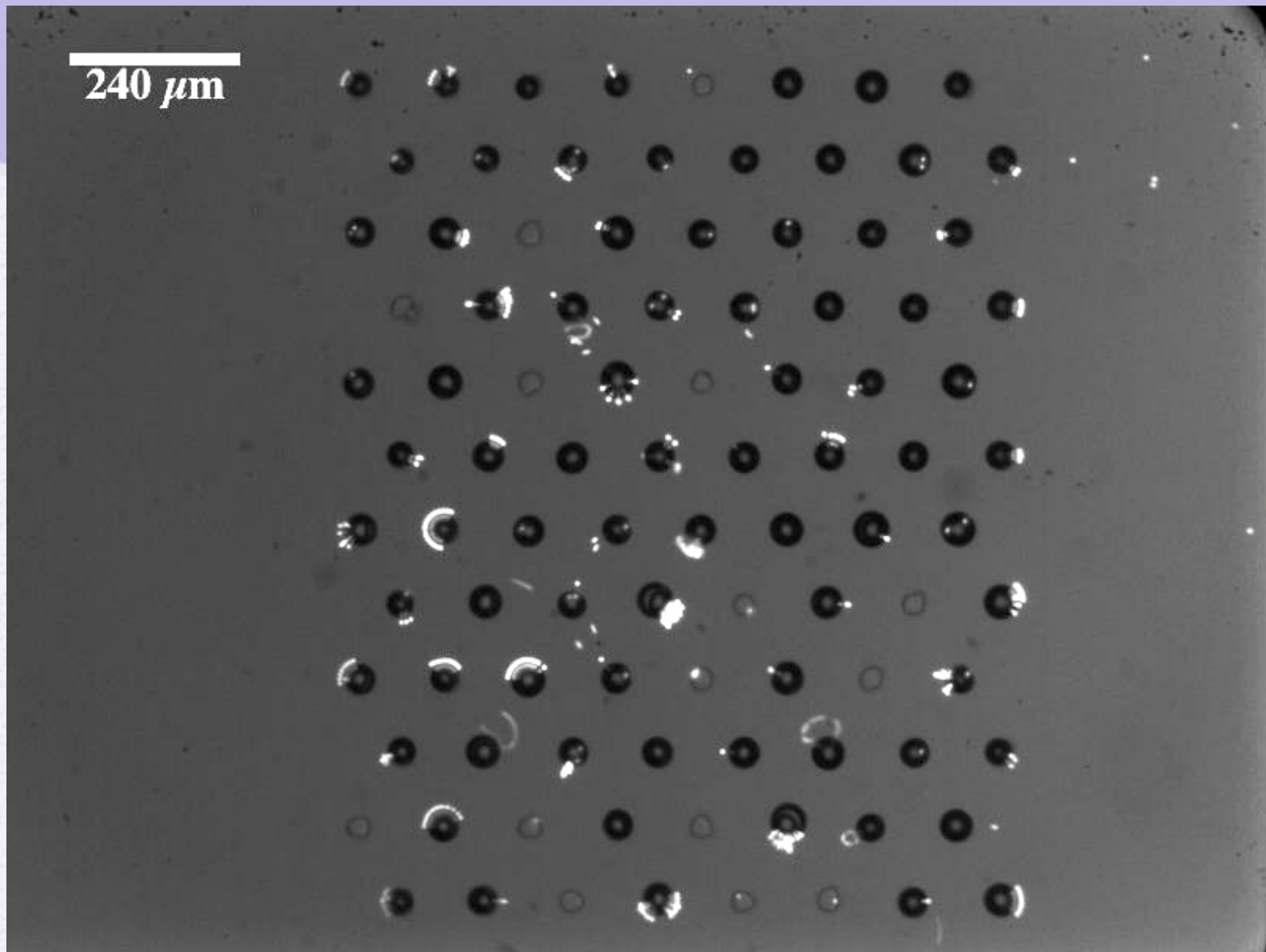
Neighbour 2

Neighbour 3

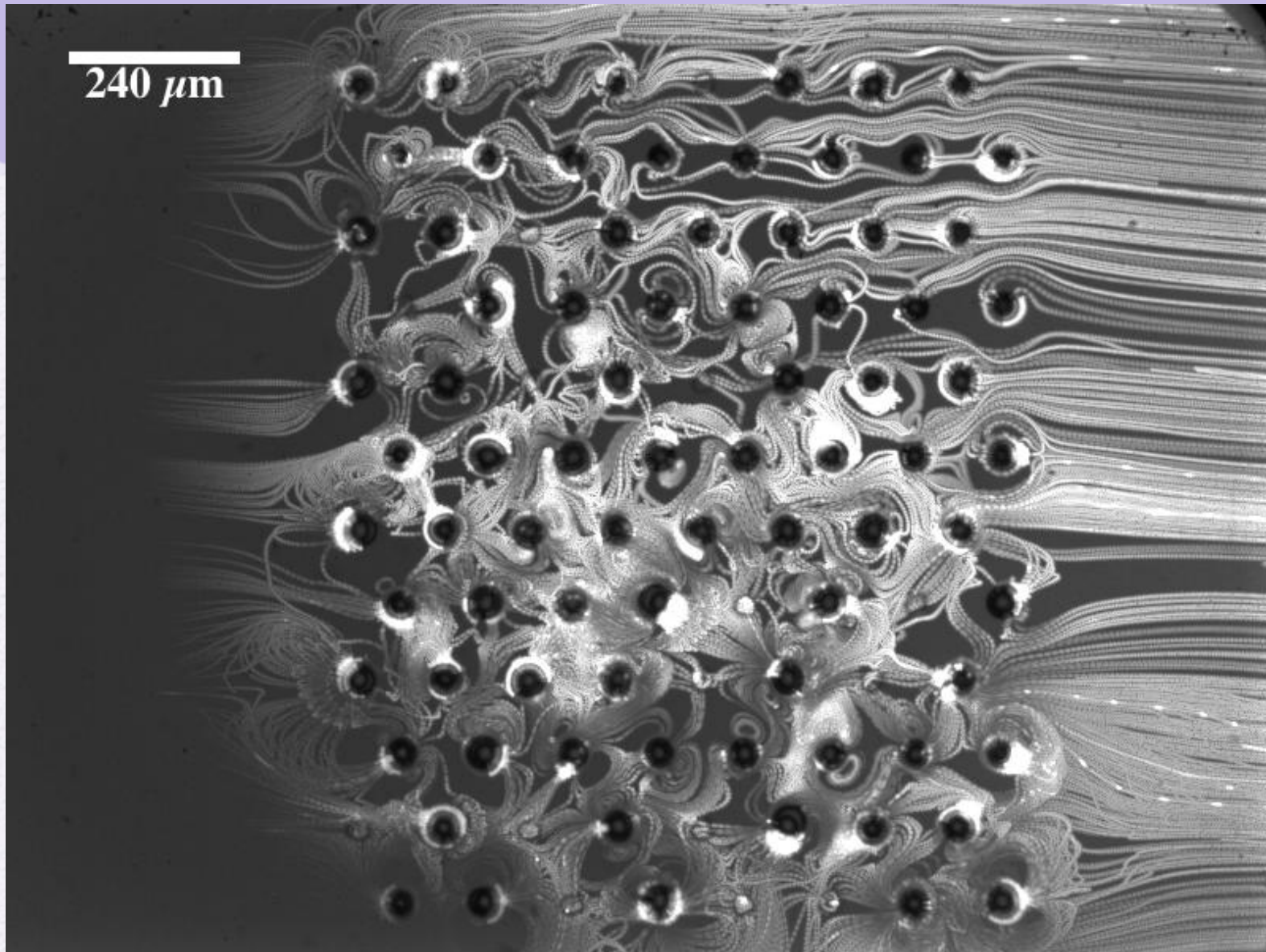
Bubble 0 under scrutiny

- Seven bubbles





- Microfluidic "pinball"



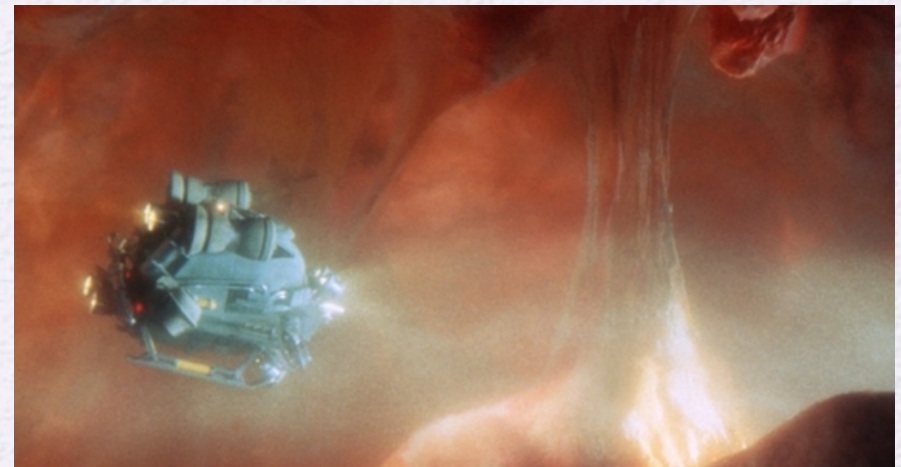
- Microfluidic "pinball"

Artificial microswimmers



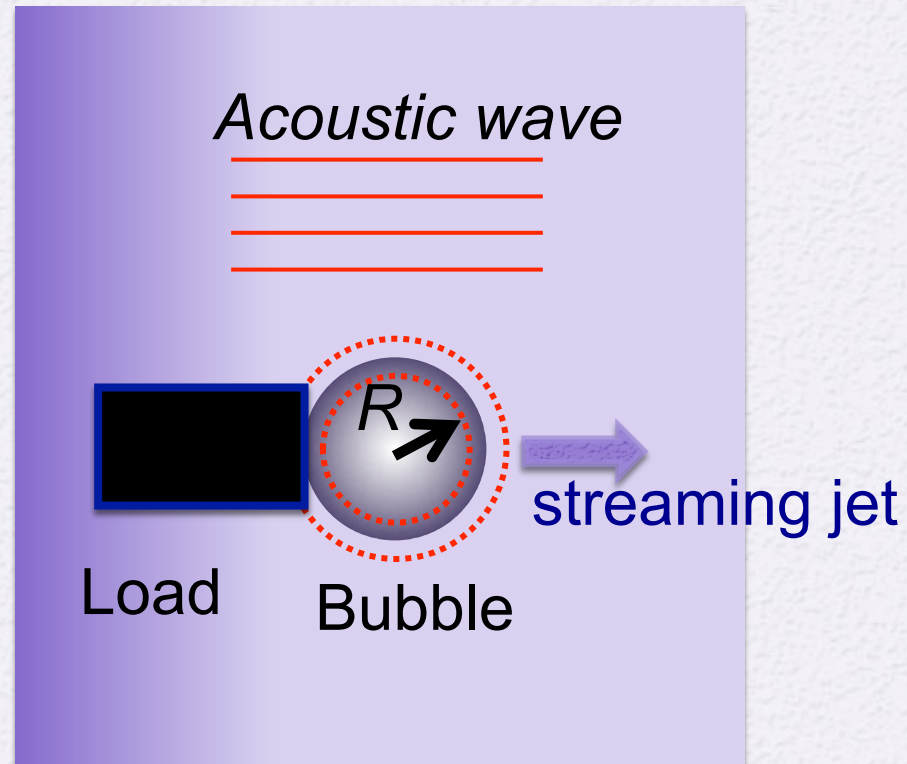
Applications

- Drug delivery in capillary vessels
- Micromixing



Innerspace, 1987

Idea: use a resonant bubble as a motor

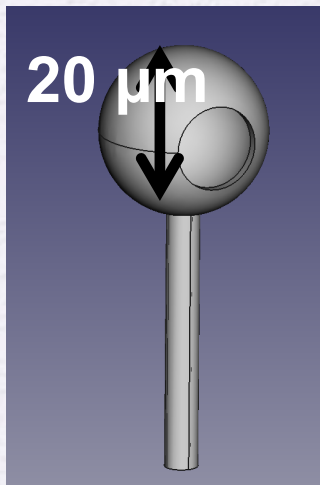


Resonance frequency
given by the Minnaert formula:

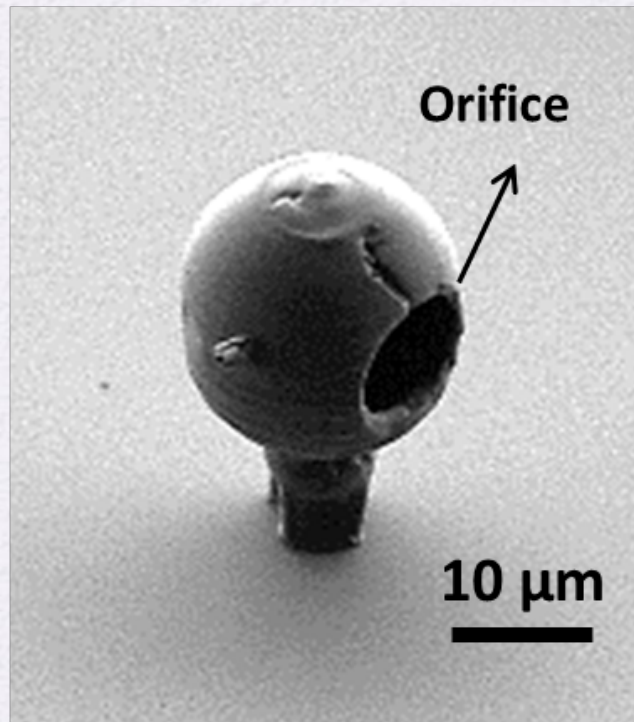
$$f \times R \simeq 3 \text{ m/s}$$

- Strong amplitude of vibration at resonance
- Issue: dissolution of the gas, especially at small scales!

How to protect a bubble from dissolving? partial encapsulation in 3D microfabricated shells

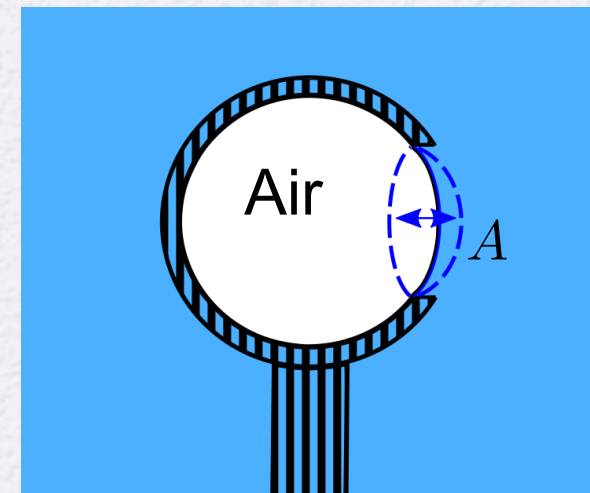


FreeCAD



Microfabrication

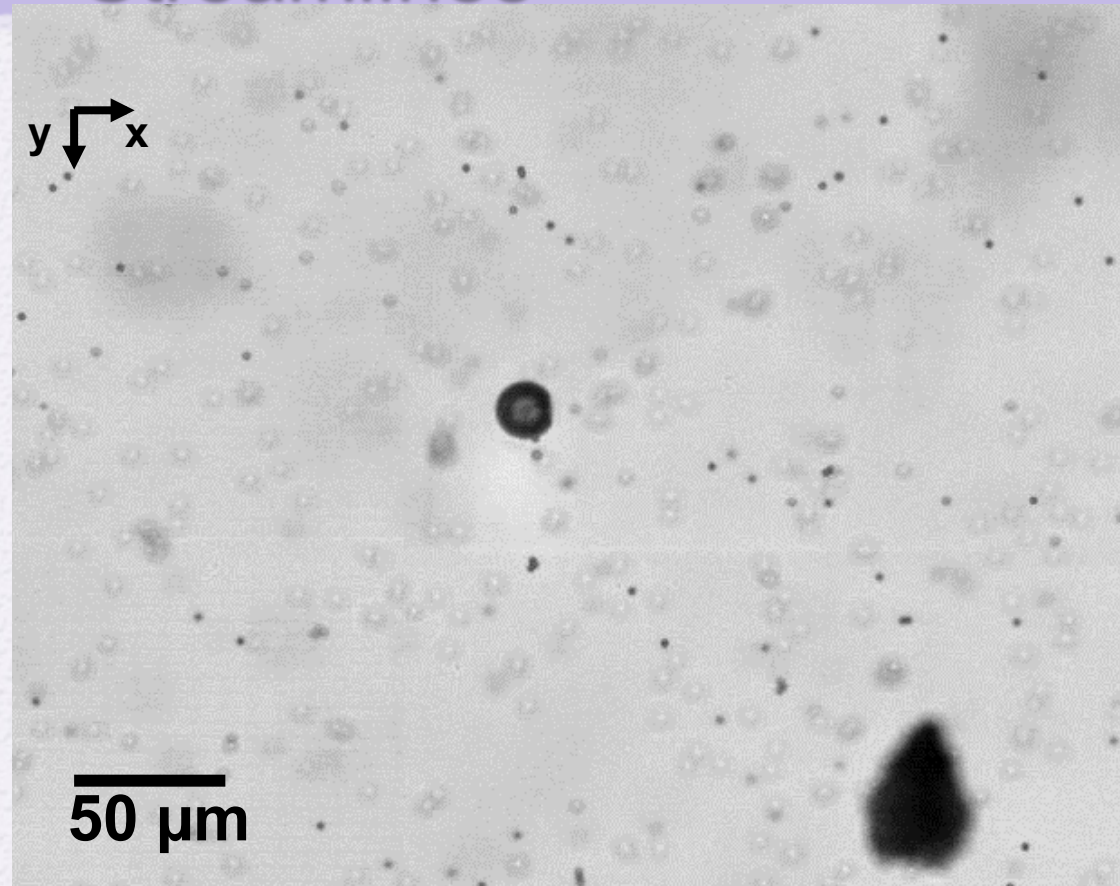
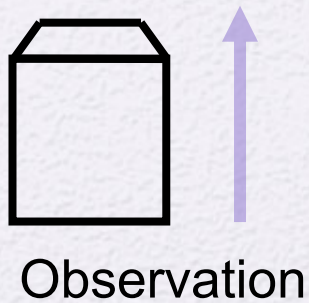
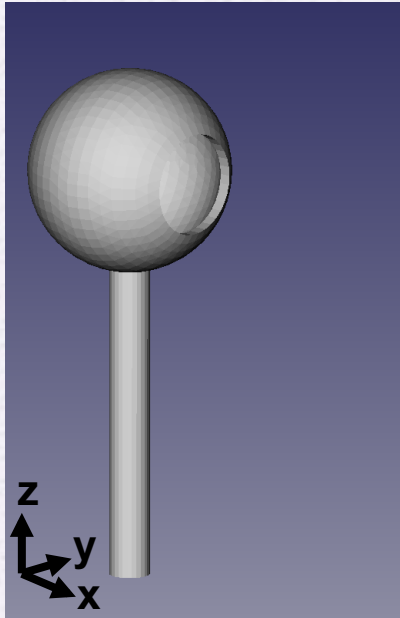
*2 photon polymerisation
0.5 μm in resolution*



In water

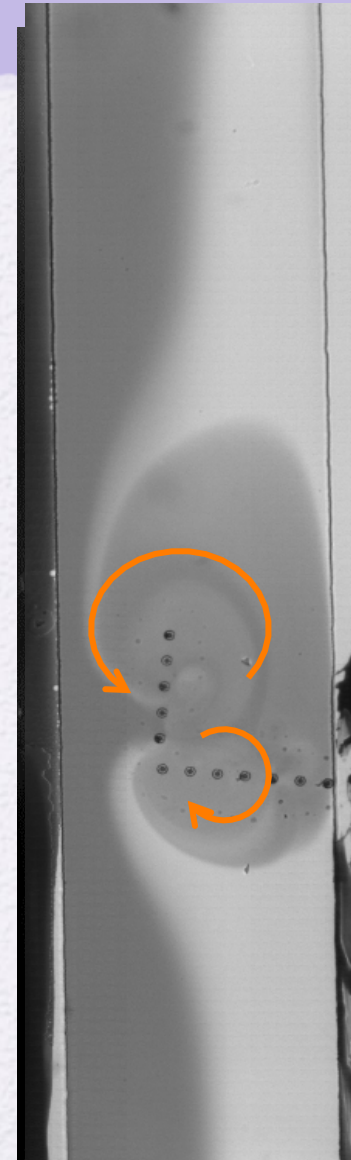
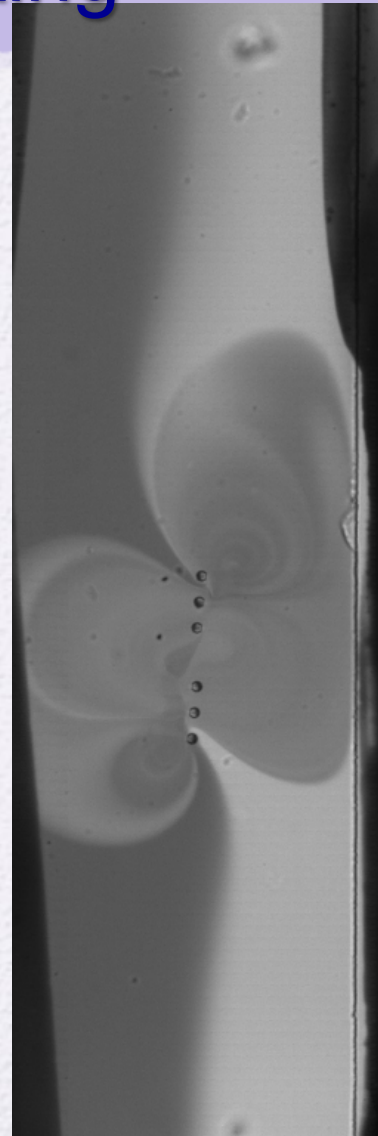
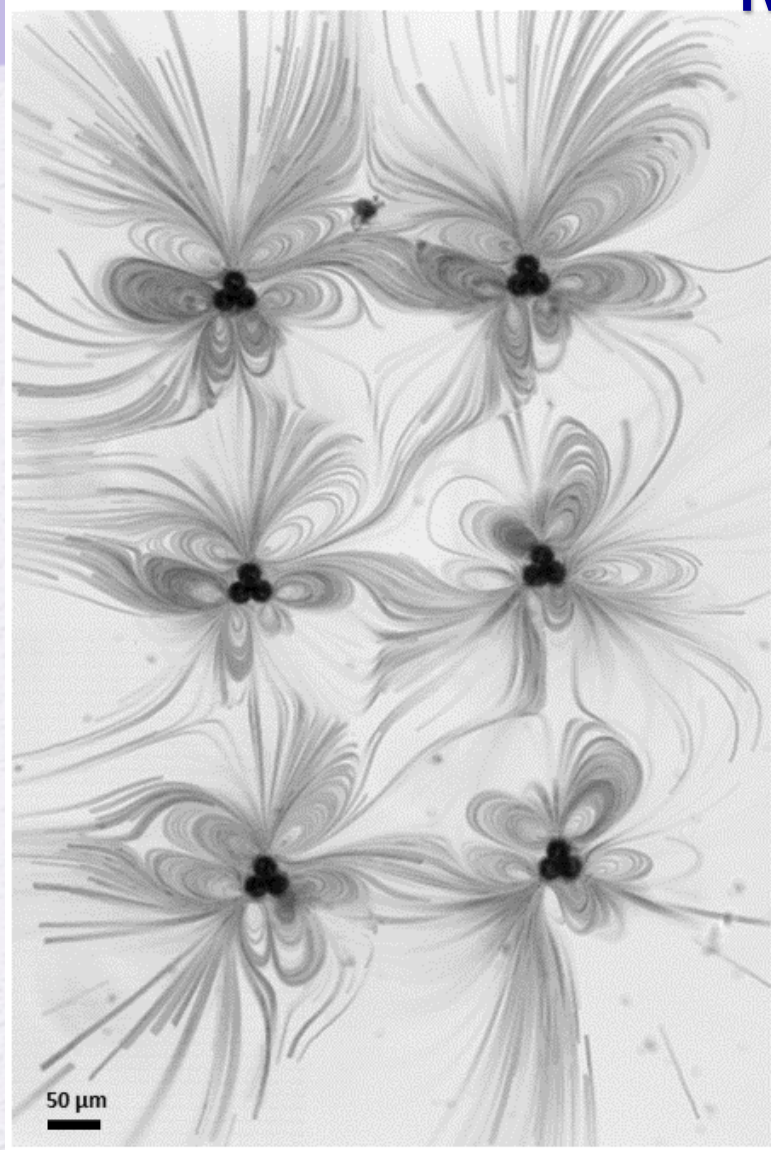
- Interface is pinned: dissolution is prevented (but not condensation)
- Large amplitudes of vibration!

Streamlines



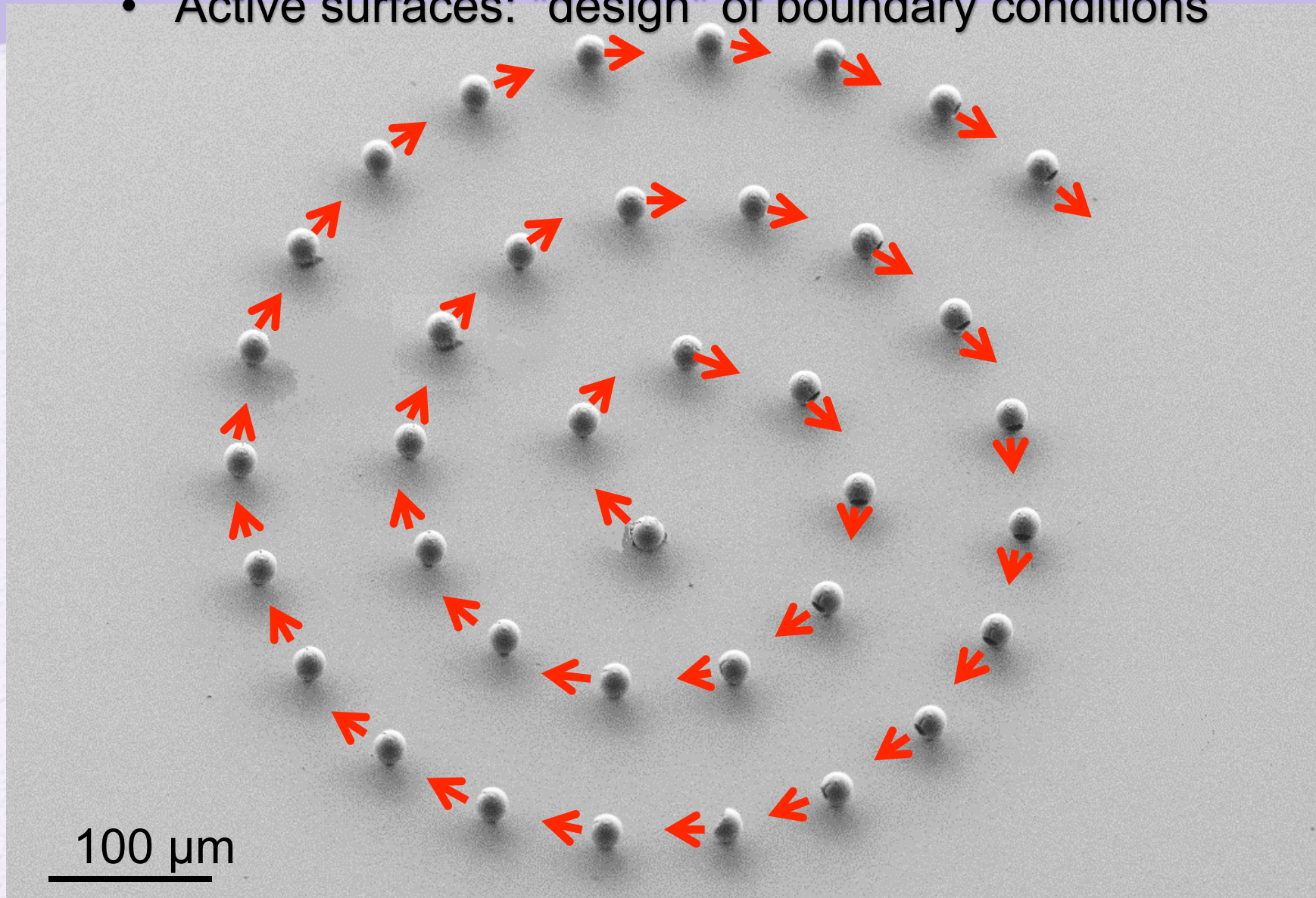
$f_{\text{transd}} = 320 \text{ kHz}$, $P_{\text{ac}} = 9.2 \text{ kPa}$,
slowed down 32x

Mixing

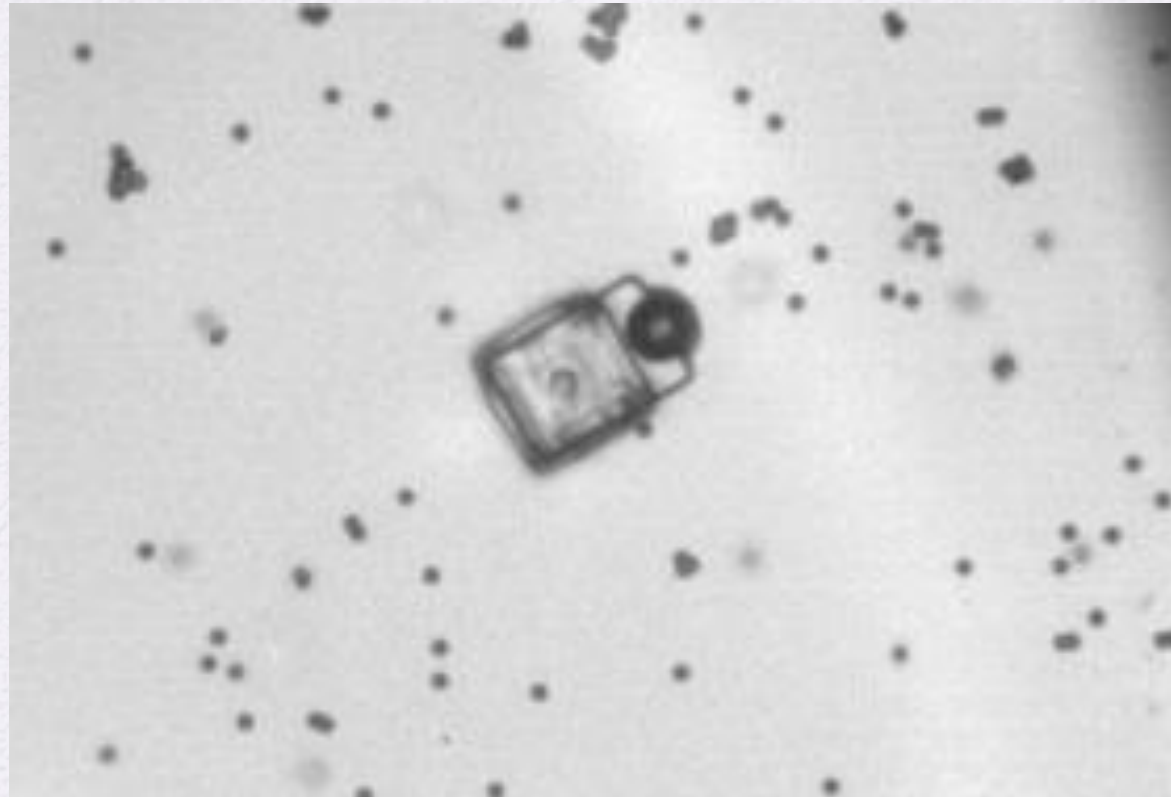


Bertin 2017

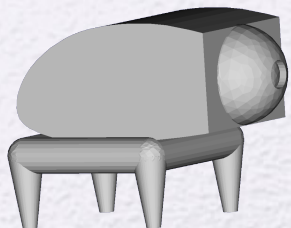
- Active surfaces: "design" of boundary conditions



Swimmer!

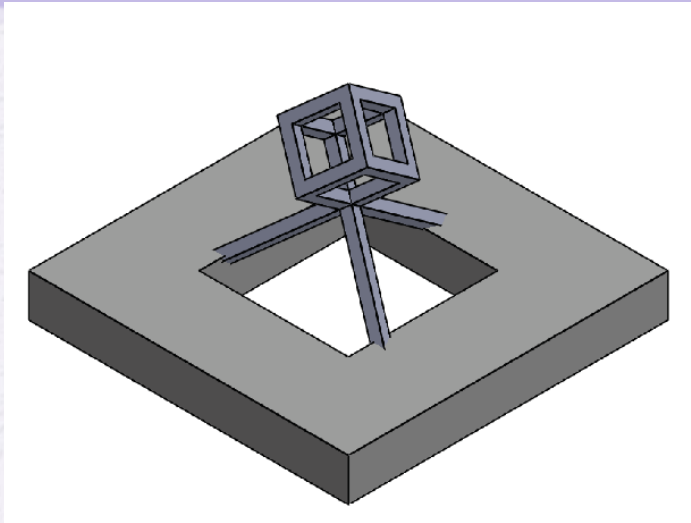


piezo

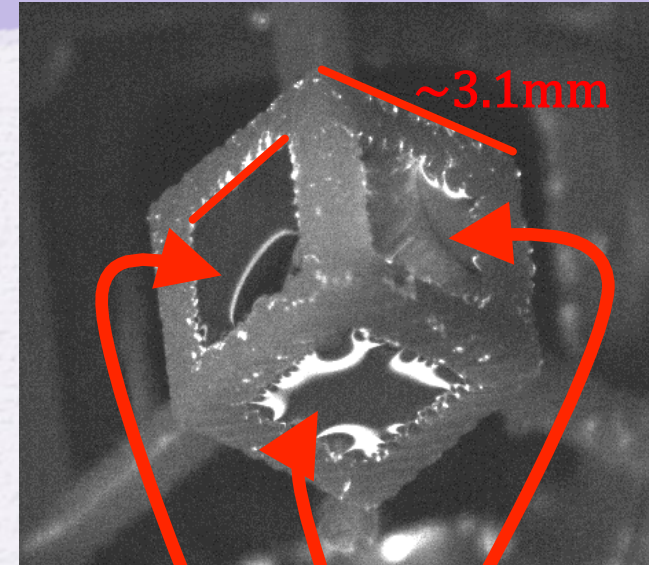


Video Jean-François Louf, $f=340$ kHz

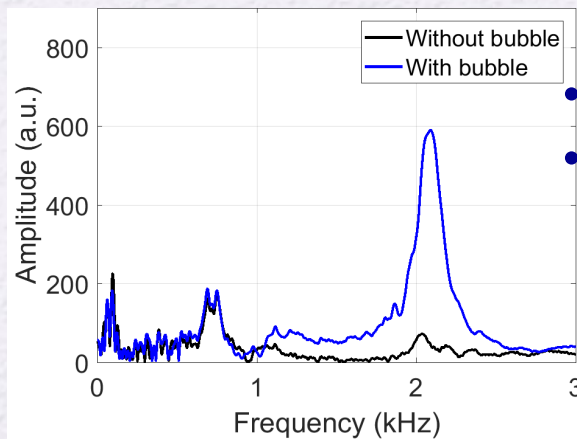
Cubic bubbles



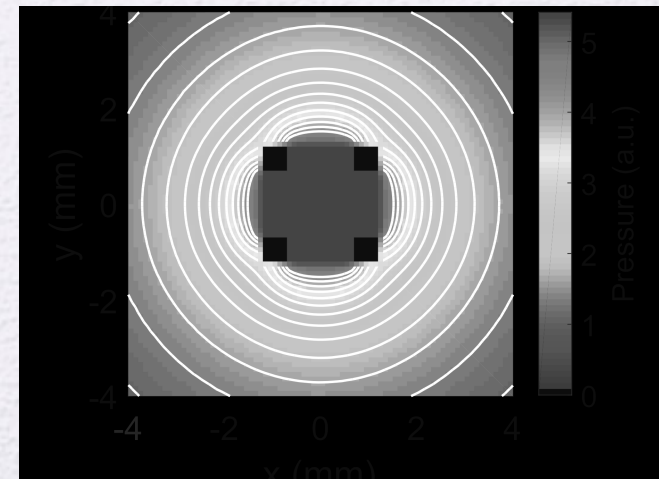
3D fabrication (SLA)



Water-air interfaces



Stable
Good resonator!



Conclusion

- Acoustics: **easy to implement**
- Strong radiation forces on objects:
 - I/c when larger than λ
 - $I/c (R/\lambda)^4$ for objects smaller
 - $I/c (R/\lambda)^1$ for standing waves
- Forces on water itself: acoustic streaming
- Care has to be taken to **focus the energy**
(use of high frequency, resonances, bubbles)

